

THE FUNCTIONAL MODEL FOR MAXIMAL DISSIPATIVE OPERATORS (TRANSLATION FORM): AN APPROACH IN THE SPIRIT OF OPERATOR KNOTS

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ABSTRACT. In this article we develop a functional model for a general maximal dissipative operator. We construct the selfadjoint dilation of such operators. Unlike previous functional models, our model is given explicitly in terms of parameters of the original operator, making it more useful in concrete applications.

For our construction we introduce an abstract framework for working with a maximal dissipative operator and its anti-dissipative adjoint and make use of the Štraus characteristic function in our setting. Explicit formulae are given for the selfadjoint dilation, its resolvent, a core and the completely non-selfadjoint subspace; minimality of the dilation is shown. The abstract theory is illustrated by the example of a Schrödinger operator on a half-line with dissipative potential, and boundary condition and connections to existing theory are discussed.

1. INTRODUCTION

In recent years, the spectral and scattering properties of non-selfadjoint problems have become a subject of much mathematical and physical interest. This is the natural setting for many important problems in physics including dissipative problems (where the system loses energy), problems in hydrodynamics and the study of metamaterials where progress has been driven in part by the development and feasibility of manufacture of novel materials with unexpected properties. Dissipation, at the atomic level, plays an essential part in many processes, see for example Milton et al. on cloaking in the presence of a superlens [29], Weder et al. [14] on plasma heating through tunneling effects in tokamaks, the work of Figotin and Welters on dissipation in composite materials [17], Cherednichenko et al. on quantum graphs using the functional model [8] and Fröhlich et al. on scattering for the Lindblad equations [16] where dissipative methods were used.

Mathematically these problems pose a challenge, as apart from rather exceptional cases, the well-developed methods used to examine the spectrum of selfadjoint problems are not applicable. According to Mark Krein the spectral theorem in the selfadjoint case highlights the relationship between the spectral analysis of the operator and the geometry of the Hilbert space; in contrast, in the spectral analysis of non-selfadjoint operators this geometric relationship plays a much reduced role and is replaced by complex analysis. A tool more appropriate to analyse the spectrum of non-selfadjoint operators has to be used; such a tool is the functional model. This reduces the spectral analysis of a non-selfadjoint operator to a problem in complex analysis: the canonical factorisation of the characteristic function as an analytic operator-valued function in the upper half-plane (M. Livšić theorem). The functional model provides a systematic approach to studying the spectral and scattering theory of non-selfadjoint problems with wide applicability.

Following pioneering work by Livšić on characteristic functions and the so-called triangular model [25, 26], functional models were introduced for contractions by Sz. Nagy and Foias (see [44, 32] and references therein) to analyse the structure of contractions and relations between an operator, its spectrum and its characteristic function, and simultaneously, in a different form, by de Branges [10]. Since then functional models have been developed further including a very useful symmetric version of the Sz.-Nagy-Foias model due to Pavlov [35]. They have been used to obtain many results in mathematical physics and in spectral analysis with applications to problems such as Schrödinger operators with complex potentials and non-selfadjoint boundary conditions, and stochastic quantum dynamics. Pavlov's work on quantum switches [38] and Naboko and Romanov's work on time asymptotics for the Boltzmann operator [31] have relied heavily on it. The best known application is Lax-Phillips scattering theory which corresponds to a special case of the Sz. Nagy-Foias functional model when the characteristic function of the operator to be studied is an inner function (this excludes the possibility of absolutely continuous spectrum). Functional models can be used to find conditions for the existence and completeness for wave operators in scattering theory, the scattering matrix and spectral shift function and give explicit formulae for them in the framework of the model (see [30]). Moreover, the functional model has applications in inverse scattering theory and can help

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provide information on which part of the operator can be reconstructed from measurements and which part can be cloaked, see e.g. [24].

Constructions of functional models often rely on abstract results, making it hard to apply the results to specific examples. Based on Pavlov's explicit construction of the functional model for Schrödinger operators in [37], the functional model for the situation of a family of additive perturbations to a selfadjoint operator was developed in [30]. The approach has several concrete advantages: It gives explicit formulae for all expressions arising in the model and, by developing the spectral form of the functional model, it importantly also allows the study of non-dissipative operators within the framework of the functional model. In further work, Ryzhov has developed a functional model for the case when the perturbation is only in the boundary conditions [39]. Recently, a functional model has been used to study transmission problems for PDEs [9].

In this paper, we consider a general maximal dissipative operator and develop the so-called 'translation form' of the functional model. This consists of constructing a selfadjoint dilation of the maximal dissipative operator by adding two so-called channels. The channels model incoming and outgoing effects and the dilation acts in them as a first order differential operator. Kudryashov [22] presented an explicit form of the selfadjoint dilation of a general maximal dissipative operator, announcing only the result without any details of the proof. Later, Ryzhov [40] reconstructed this form of the selfadjoint dilation and gave a complete proof. In both cases the explicit construction of the dilation is purely abstract and contains the square roots of some complicated operators which can only be explicitly calculated in a useful form in the case of a rank one perturbation of a selfadjoint operator.

Our goal in this paper is to present an alternative construction of the selfadjoint dilation of a maximal dissipative operator using an abstract construction in the spirit of operator knots [5]. One of the main ingredients of that theory is a Lagrange identity which we here extend from the case of bounded operators to our more general setting. This relates functions in the domain of the maximal dissipative operator via so-called Γ -operators to the non-selfadjoint part of the operator. The approach is also inspired by the theory of boundary triples (see, e.g. [7, 11, 18, 20]), where the Lagrange identity is replaced by a Green formula which relates functions in the domain of an operator with their boundary data. In our view, the flexibility of the choice of the Γ -operators in the Lagrange identity gives a serious advantage to our construction. It will enable us to choose the Γ -operators so that expressions arising in the dilation can be given explicitly in terms of model parameters such as coefficients or boundary conditions of the maximal dissipative operator. This will allow us to obtain spectral results in explicit terms of these model parameters which is not possible in the constructions in [22] and [40].

This paper is structured in the following way: Section 2 gives a brief introduction to the main facts about dissipative operators and selfadjoint dilations and an overview of the historical development. The Sections 3-4 discuss background material and introduce concepts and notation required. Although the ideas in these sections may not be completely new, many of them can already be found in the work of Štraus for special cases, to the best of our knowledge the presentation in full generality here is new and additionally serves to make our presentation self-contained. In Section 3 we introduce an abstract framework for working with a maximally dissipative operator and its anti-dissipative adjoint. The Štraus characteristic function is introduced in our framework in Section 4 and its relation to the Sz.-Nagy-Foias characteristic function for contractions is discussed. The Sections 5-8 cover the main new material of this paper. Section 5 introduces our formula for the selfadjoint dilation and its domain. This is illustrated for two examples, including a dissipative Schrödinger operator on a half-line, in Section 6, showing the advantages of our explicit construction in the spirit of operator knots. Properties of the dilation are shown in Section 7, in particular, we calculate the resolvent and show minimality of the dilation. As a by-product we present a relatively simple expression for the dilation on a core. We also discuss complete non-selfadjointness of the dilation, giving an explicit formula for the completely non-selfadjoint subspace. Finally, in Section 8, we discuss connections of our construction with the Kudryashov/Ryzhov model and look more closely at the characteristic function in the case of a symmetric minimal operator, even in the case of differing deficiency indices.

The paper develops the translation form of the functional model. Using it, we will develop the so-called spectral form of the functional model in a forthcoming paper. In the spectral form, the operator is represented as the operator of multiplication by an independent variable in some auxiliary vector-valued function space; in our version the connection between this auxiliary space and the original Hilbert space will be explicit.

Throughout the paper we use the following notation: For a complex number $z \in \mathbb{C}$, let $\Im z$ denote its imaginary part and $\mathbb{C}^+ = \{z \in \mathbb{C} : \Im z > 0\}$, $\mathbb{C}^- = \{z \in \mathbb{C} : \Im z < 0\}$. The positive half-line will be denoted by \mathbb{R}^+ . For an operator A in a Hilbert space H , we denote its range by $\text{Ran } A$, its kernel by $\ker(A)$, its adjoint by A^* and its spectrum and resolvent set by $\sigma(A)$ and $\rho(A)$, respectively. The set of bounded linear operators in H is denoted $B(H)$. The Lebesgue space of square-integrable functions on the half-line is $L^2(\mathbb{R}^+)$, while $H^s(\mathbb{R}^+)$ denotes the usual Sobolev space of order s ; $H_0^s(\mathbb{R}^+)$ denotes the closure in H^s -norm of the smooth, compactly supported functions on the half-line.

2. BASIC PROPERTIES OF DISSIPATIVE OPERATORS

This section reviews some classical results on dissipative operators. For more on the subject, we refer the reader to [19, 24, 44].

Definition 2.1. *A densely defined linear operator A with domain $D(A)$ in H is called dissipative if $\Im \langle Ah, h \rangle \geq 0$ for all $h \in D(A)$. A is called anti-dissipative if $(-A)$ is dissipative.*

The next result shows that every $\lambda \in \mathbb{C}^-$ is of regular type for a dissipative operator A .

Proposition 2.2. *Let A be dissipative. For any $f \in D(A)$ and $\lambda \in \mathbb{C}^-$ we have*

$$\|(A - \lambda)f\| \geq |\Im \lambda| \|f\|.$$

One would like to additionally have invertibility of $A - \lambda$, i.e. $\text{Ran}(A - \lambda) = H$. The following theorem guarantees that every densely defined dissipative operator has an extension with this property:

Proposition 2.3 (R. Phillips). *For any densely defined dissipative operator A there exists at least one dissipative extension \tilde{A} , i.e. $D(A) \subseteq D(\tilde{A})$ and $\tilde{A}|_{D(A)} = A$, such that $\text{Ran}(\tilde{A} - \lambda) = H$ for all $\lambda \in \mathbb{C}^-$.*

Definition 2.4. *Dissipative operators which have no non-trivial dissipative extensions are called maximal dissipative operators (MDO).*

MDOs are characterised by the existence of a bounded resolvent $(A - \lambda)^{-1}$ on the whole of H for $\lambda \in \mathbb{C}^-$, i.e. MDOs have no spectrum in the lower half plane. A dissipative operator may have several maximal dissipative extensions.

There is a bijective map between the class of MDOs and contractions which is given by the Cayley transform: For any MDO, $(A + i)^{-1} \in B(H)$. Define

$$(2.1) \quad T = I - 2i(A + i)^{-1} = (A - iI)(A + iI)^{-1}.$$

This is an operator version of the Möbius transform.

Properties of the Cayley transform:

- (1) Let T be the Cayley transform of an MDO. Then $D(T) = H$ and $\|T\| \leq 1$.
- (2) The Cayley transform is a one-to-one map from the class of MDOs in H onto the class of all contractions satisfying the extra condition $1 \notin \sigma_p(T)$. Here, $\sigma_p(T)$ denotes the set of eigenvalues of T . The condition $1 \notin \sigma_p(T)$ is equivalent to $\text{Ran}(T - I)$ being dense in H .
- (3) $\lambda \in \sigma(A) \subseteq \overline{\mathbb{C}^+}$ iff $(\lambda - i)/(\lambda + i) \in \sigma(T)$.

Remark 2.5. *At first glance the Cayley transform looks like a very convenient tool, replacing the complicated class of unbounded MDO by the class of contractions. However, for a particular MDO A there is rarely sufficient information on $(A + i)^{-1}$ to explicitly obtain T . Therefore, the importance of the Cayley transform is often purely theoretical.*

The real eigenvalues of an MDO exhibit the same behaviour as the eigenvalues of selfadjoint operators.

Proposition 2.6 (Sz.Nagy). *Let A be an MDO. Then the eigenvectors corresponding to real eigenvalues are orthogonal to all eigenvectors corresponding to different eigenvalues (real or complex). Moreover, the subspace spanned by all eigenvectors corresponding to real eigenvalues belongs to the selfadjoint subspace H_1 in the Langer decomposition (see Proposition 2.10).*

Remark 2.7. *Note that for an MDO A there cannot be any root vectors corresponding to real spectrum. Just as in the well-known case of matrices, this follows from the resolvent estimate $\|(A - \lambda I)^{-1}\| \leq (|\Im \lambda|)^{-1}$ for $\lambda \in \mathbb{C}^-$.*

The proposition says that our operator consists of a part (corresponding to the set of eigenvectors of the real point spectrum) which looks like a selfadjoint operator and a remaining part. It seems reasonable to try to study the two parts separately. This idea leads to the introduction of the notion of completely non-selfadjoint operators (corresponding to the remaining part of the operator).

Definition 2.8. *Let A be an operator on a Hilbert space H , $H_1 \subseteq H$ a subspace and P_{H_1} the orthogonal projection of H onto H_1 . The subspace H_1 is invariant with respect to A if $P_{H_1}D(A) \subseteq D(A)$ and $AP_{H_1}h \in H_1$ for all $h \in D(A)$. It is a reducing subspace for A if both H_1 and $H \ominus H_1$ are invariant with respect to A .*

Definition 2.9. *Let A be an MDO. A is completely non-selfadjoint (cns) if there exists no reducing subspace $H_1 \subseteq H$ such that $A|_{H_1}$ is selfadjoint.*

The following result gives an explicit formula for the completely non-selfadjoint part of the operator. In the case of relatively bounded imaginary part the formula is simple. For more general situations the formula involves operators Δ and Δ_* which are regularisations of the (possibly non-existing) imaginary part of the operator. In our setting, we will determine an explicit formula for the completely non-selfadjoint part of an MDO in Theorem 7.6.

Proposition 2.10. (*Langer decomposition, see [23, 30]*). *Let A be an MDO. Then there exists a unique decomposition of $H = H_1 \oplus H_2$ into an orthogonal sum of two reducing subspaces for A such that $A|_{H_1}$ is selfadjoint in H_1 and $A|_{H_2}$ is completely non-selfadjoint in H_2 .*

Define

$$(2.2) \quad \Delta = I - T^*T = 2i \left[(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1}(A + i)^{-1} \right],$$

$$(2.3) \quad \Delta_* = I - TT^* = 2i \left[(A + i)^{-1} - (A^* - i)^{-1} + 2i(A + i)^{-1}(A^* - i)^{-1} \right]$$

and set

$$\mathcal{M} := \text{Ran}(\Delta) + \text{Ran}(\Delta_*) \subseteq H.$$

Then the completely non-selfadjoint part H_2 is given by

$$H_2 = \text{clos} \left(\text{Span}_{\Im \lambda < 0} \{ (A - \lambda)^{-1} \mathcal{M} \} + \text{Span}_{\Im \lambda > 0} \{ (A^* - \lambda)^{-1} \mathcal{M} \} \right).$$

If A has relatively bounded imaginary part, i.e. $A = L + iV$ with $L = L^*$, $V \geq 0$, V relatively L -bounded, then there is a simple explicit expression for the completely non-selfadjoint part H_2 :

$$H_2 = \text{clos} \left(\text{Span}_{\Im \lambda \neq 0} \{ (L - \lambda)^{-1} \text{Ran } V \} \right) = \text{clos} \left(\text{Span}_{\lambda \notin (\sigma(A) \cup \mathbb{R})} \{ (A - \lambda)^{-1} \text{Ran } V \} \right),$$

i.e. H_2 is generated by the range of the imaginary part V developed by the resolvent of the operator A or its real part L . Moreover, $A|_{H_1} = L|_{H_1}$.

In systems theory, MDOs are used to describe systems with a loss of energy, while Hermitian operators describe systems with energy conservation. This naturally leads to the idea of including a dissipative system in a larger conservative one, taking into account ‘where’ the energy is leaking to. The mathematical realization of this idea is due to the Hungarian mathematician B. Sz.-Nagy in the late 50ies, but its roots go back to earlier papers by M. Naimark. Actually, Sz.-Nagy worked with contractions rather than MDOs. However the two formulations are equivalent via the Cayley transform. See [44, Theorem I.4.1].

Proposition 2.11 (Sz.-Nagy). *For any MDO A on a Hilbert space H there exists a selfadjoint operator \mathcal{L} on a Hilbert space $\mathcal{H} \supseteq H$ such that*

$$(2.4) \quad e^{itA} = P_H e^{it\mathcal{L}}|_H, \quad t \geq 0 \quad \text{or equivalently} \quad (A - \lambda)^{-1} = P_H (\mathcal{L} - \lambda)^{-1}|_H, \quad \lambda \in \mathbb{C}^-.$$

The operator \mathcal{L} is called a selfadjoint dilation of A .

Definition 2.12. *A dilation is minimal if it contains no non-trivial reducing part which is itself a selfadjoint dilation of A .*

The minimal selfadjoint dilation of an arbitrary MDO A will be the sum of the selfadjoint part of A and the minimal selfadjoint dilation of the completely non-selfadjoint part. Any completely non-selfadjoint operator has a minimal selfadjoint dilation. The following result is due to Sz.-Nagy and Foiaş [43], see also [44, Theorem II.6.4].

Proposition 2.13. *The minimal selfadjoint dilation of a completely non-selfadjoint MDO A always has pure absolutely continuous spectrum covering the whole real line, in particular $d\langle E_\lambda^\mathcal{L} h, h \rangle$ is an absolutely continuous measure for any $h \in \mathcal{H}$, it is equivalent to the Lebesgue measure on \mathbb{R} for any $h \in H \setminus \{0\}$ and generically for $h \in \mathcal{H}$. Here, $E_\lambda^\mathcal{L}$ is the spectral resolution of the selfadjoint dilation \mathcal{L} .*

Corollary 2.14. *Let A be an MDO such that the spectrum of its minimal dilation does not cover the whole real line. Then $A = A^*$.*

3. THE LAGRANGE IDENTITY

Boundary triples are a way of naturally associating ‘boundary operators’ with an adjoint pair of operators. In the abstract setting, Weyl functions can be introduced and many questions e.g. concerning the extension theory of operators can be investigated in the framework, see e.g. [7, 11] for details. We now discuss a similar abstract framework for a maximally dissipative operator and its anti-dissipative adjoint which allows us to introduce Γ -operators associated with the imaginary part of the operator A . For the case of bounded operators this goes back to the work of the Odessa school on operator knots [5], see also [41].

Lemma 3.1. *Let A be a maximally dissipative operator on a Hilbert space H . Then there exists a Hilbert space E and an operator $\Gamma : D(A) \rightarrow E$ which is bounded in the graph norm of A , has dense range in E and such that for all $u, v \in D(A)$ we have*

$$(3.1) \quad \langle Au, v \rangle_H - \langle u, Av \rangle_H = i \langle \Gamma u, \Gamma v \rangle_E.$$

Similarly, there exists a Hilbert space E_ and an operator $\Gamma_* : D(A^*) \rightarrow E_*$ which is bounded in the graph norm, has dense range in E_* and such that for all $u, v \in D(A^*)$ we have*

$$(3.2) \quad \langle A^*u, v \rangle_H - \langle u, A^*v \rangle_H = -i \langle \Gamma_*u, \Gamma_*v \rangle_{E_*}.$$

Proof. Define the sesquilinear form

$$a[u, v] := \frac{1}{i} (\langle Au, v \rangle_H - \langle u, Av \rangle_H) \quad \text{for } u, v \in D(A).$$

Since A is dissipative, a is positive. Moreover, since A is maximal dissipative, $(A + i)^{-1}$ exists and we can define another positive sesquilinear form

$$b[f, g] := a[(A + i)^{-1}f, (A + i)^{-1}g] \quad \text{for } f, g \in H.$$

Note that for $u \in D(A)$,

$$2\Im \langle Au, u \rangle \leq \|Au\|^2 + \|u\|^2 \quad \Rightarrow \quad \Im \langle Au, u \rangle \leq \frac{1}{4} (\|Au\|^2 + 2\Im \langle Au, u \rangle + \|u\|^2).$$

Then, with $u = (A + i)^{-1}f$, for the quadratic form we have

$$|b[f, f]| = 2|\Im \langle Au, u \rangle_H| \leq \frac{1}{2} (\|Au\|^2 + 2\Im \langle Au, u \rangle + \|u\|^2) = \frac{1}{2} \|(A + i)u\|^2 = \frac{1}{2} \|f\|^2.$$

Therefore, by the Riesz representation theorem [33], there exists a non-negative bounded operator $F_b : H \rightarrow H$ with $\|F_b\| \leq 1/\sqrt{2}$ such that

$$b[f, g] = \langle F_b f, g \rangle_H = \left\langle \sqrt{F_b} f, \sqrt{F_b} g \right\rangle_H, \quad f, g \in H$$

and so

$$a[u, v] = \left\langle \sqrt{F_b}(A + i)u, \sqrt{F_b}(A + i)v \right\rangle_H, \quad u, v \in D(A).$$

Let $E = \overline{\text{Ran}(\sqrt{F_b})} = \overline{\text{Ran}(F_b)}$ equipped with the scalar product from H and set $\Gamma = \sqrt{F_b}(A + i) : D(A) \rightarrow E$. Then E and Γ have the required properties and (3.1) holds.

To obtain E_* and Γ_* , repeat the same construction for the maximal dissipative operator $-A^*$. \square

Remark 3.2. *In general, E and E_* may be of different dimensions, as can be seen in the examples below. However, in the special case of bounded imaginary part of A , we can always choose $E = E_*$ and $\Gamma = \Gamma_*$.*

The next lemma (see also [41]) shows that the operator Γ is uniquely determined up to unitary transformations.

Lemma 3.3. *Assume that for any $u, v \in D(A)$ we have*

$$\langle Au, v \rangle_H - \langle u, Av \rangle_H = i \langle \Gamma u, \Gamma v \rangle_E = i \langle \tilde{\Gamma} u, \tilde{\Gamma} v \rangle_{\tilde{E}},$$

where $\Gamma : D(A) \rightarrow E$ and $\tilde{\Gamma} : D(A) \rightarrow \tilde{E}$ are linear maps with dense range in some auxiliary Hilbert spaces E and \tilde{E} , respectively.

Then there exists a unitary map $V : \tilde{E} \rightarrow E$ such that $\Gamma u = V\tilde{\Gamma}u$ for all $u \in D(A)$.

Proof. We have that $\langle \Gamma u, \Gamma v \rangle_E = \langle \tilde{\Gamma} u, \tilde{\Gamma} v \rangle_{\tilde{E}}$ for any $u, v \in D(A)$. Assume that $h \in \text{Ran}(\Gamma)$, i.e. $h = \Gamma u$ for some $u \in D(A)$. Assume that $h = \Gamma u'$ for some other $u' \in D(A)$. Then

$$\langle \tilde{\Gamma} u, \tilde{\Gamma} v \rangle_{\tilde{E}} = \langle \Gamma u, \Gamma v \rangle_E = \langle \Gamma u', \Gamma v \rangle_E = \langle \tilde{\Gamma} u', \tilde{\Gamma} v \rangle_{\tilde{E}}.$$

Since $\text{Ran}(\tilde{\Gamma})$ is dense in \tilde{E} , we get $\tilde{\Gamma}u = \tilde{\Gamma}u'$. Therefore the map $V : \text{Ran}(\Gamma) \rightarrow \text{Ran}(\tilde{\Gamma})$ given by $h = \Gamma u \mapsto \tilde{h} = \tilde{\Gamma}u$ is one-to-one from $\text{Ran}(\Gamma)$ onto $\text{Ran}(\tilde{\Gamma})$ (by symmetry of Γ and $\tilde{\Gamma}$). As Γ and $\tilde{\Gamma}$ are linear, also V is linear.

Setting $u = v$ we get

$$\|h\|^2 = \|\Gamma u\|^2 = \|\tilde{\Gamma}u\|^2 = \|\tilde{h}\|^2.$$

Thus V is a unitary map from $\text{Ran}(\tilde{\Gamma})$ onto $\text{Ran}(\Gamma)$. Its closure is a unitary operator V from $\tilde{E} = \overline{\text{Ran}(\tilde{\Gamma})}$ onto $E = \overline{\text{Ran}(\Gamma)}$ such that

$$\Gamma u = V\tilde{\Gamma}u, \quad \text{for all } u \in D(A),$$

as required. \square

Despite the lemma formally showing the ‘uniqueness’ of Γ , its content is purely abstract and of little consequence in applications to particular examples. In most concrete applications, the construction using the square root used in the proof of Lemma 3.1 does not lead to explicit formulae for the operators Γ and Γ_* . However, in the following we will not make use of this construction of Γ and Γ_* . The theory will instead be valid whenever the identity (3.1) holds and Γ and E have the properties stated in the lemma. This is very much in the spirit of the boundary triples approach mentioned above. In our case, we use an abstract Lagrange identity, instead of an abstract Green identity.

Besides the choice of E, E_*, Γ and Γ_* made in the lemma, this approach allows us the freedom of choosing the operators Γ, Γ_* as two versions of the ‘roots’ of the ‘imaginary part of A ’. Even in cases when the roots do not exist, this approach allows us to give meaning to the ‘roots’, and in cases when the roots exist, it enables us to choose an alternative, simpler version of the ‘root’. In particular examples, this allows us to choose factorisations (3.1) and (3.2) which depend explicitly on parameters of the problem (such as coefficients of a differential expression). Already for the case of a rank two dissipative perturbation of a selfadjoint operator the square root is not explicit, while it is easy to find the ‘correct’ choice of Γ in the Lagrange identity. This is illustrated in the following examples.

Example 3.4. (1) *We consider a Schrödinger operator with dissipative potential and dissipative boundary condition: On $H = L^2(\mathbb{R}^+)$, let*

$$(Af)(x) = -f''(x) + q(x)f(x),$$

where q is a measurable and bounded complex-valued function on \mathbb{R}^+ with $\Im q(x) \geq 0$ for a.e. $x \in \mathbb{R}^+$ and

$$D(A) := \{f \in H^2(\mathbb{R}^+) : f'(0) = hf(0)\}$$

where $h \in \mathbb{C}$ with $\Im(h) \geq 0$. (The two conditions on the imaginary parts of q and h are necessary and sufficient for A to be maximal dissipative.) Then for $u, v \in D(A)$, we have

$$\begin{aligned} \langle Au, v \rangle - \langle u, Av \rangle &= u'(0)\overline{v(0)} - u(0)\overline{v'(0)} + 2i \int_0^\infty \Im q(x) u(x)\overline{v(x)} dx \\ &= 2i \left(\Im h u(0)\overline{v(0)} + \int_0^\infty \Im q(x) u(x)\overline{v(x)} dx \right). \end{aligned}$$

Let $\Omega = \{x \in \mathbb{R} : \Im q(x) > 0\}$, set $E = \mathbb{C} \oplus L^2(\Omega)$ and

$$\Gamma u = \begin{pmatrix} \sqrt{2\Im h} u(0) \\ \sqrt{2\Im q} u|_\Omega \end{pmatrix}.$$

Then (3.1) holds. We remark that in this example $E_* = E$ and Γ_* acts in the same way as Γ , but has a different domain.

(2) *The next simple example shows that the boundary operators Γ and Γ_* and the spaces E and E_* can differ significantly. Let*

$$A = i \frac{d}{dx} \quad \text{with} \quad D(A) = H_0^1(\mathbb{R}^+).$$

Then it is easy to check that A , being symmetric, is a maximally dissipative operator, and

$$A^* = i \frac{d}{dx} \quad \text{with} \quad D(A^*) = H^1(\mathbb{R}^+),$$

is an anti-dissipative operator, and we can choose $\Gamma = 0$ with $E = \{0\}$ and $\Gamma_* u = u(0)$ with $E_* = \mathbb{C}$.

We conclude this section with two useful identities which follow from the Lagrange identity.

Lemma 3.5. (Abstract Green Function Identities) *For $\lambda \in \mathbb{C}^+$ and $\mu \in \mathbb{C}^-$ we have*

$$(3.3) \quad (A + \lambda)^{-1} - (A^* + \mu)^{-1} + (\lambda - \mu)(A^* + \mu)^{-1}(A + \lambda)^{-1} = -i(\Gamma(A + \bar{\mu})^{-1})^*(\Gamma(A + \lambda)^{-1})$$

and

$$(3.4) \quad (A + \lambda)^{-1} - (A^* + \mu)^{-1} + (\lambda - \mu)(A + \lambda)^{-1}(A^* + \mu)^{-1} = -i(\Gamma_*(A^* + \bar{\lambda})^{-1})^*(\Gamma_*(A^* + \mu)^{-1})$$

Proof. The first result is equivalent to

$$\begin{aligned} \langle (A + \lambda)^{-1}f, g \rangle - \langle (A^* + \mu)^{-1}f, g \rangle &= -i\langle (\Gamma(A + \lambda)^{-1})f, (\Gamma(A + \bar{\mu})^{-1})g \rangle \\ &\quad - (\lambda - \mu)\langle (A + \lambda)^{-1}f, (A + \bar{\mu})^{-1}g \rangle \end{aligned}$$

for all $f, g \in H$. Defining $f_\lambda = (A + \lambda)^{-1}f$, $g_\mu = (A + \bar{\mu})^{-1}g$, this in turn is equivalent to

$$\langle f_\lambda, (A + \bar{\mu})g_\mu \rangle - \langle (A + \lambda)f_\lambda, g_\mu \rangle = -i\langle \Gamma f_\lambda, \Gamma g_\mu \rangle - (\lambda - \mu)\langle f_\lambda, g_\mu \rangle \text{ for all } f_\lambda, g_\mu \in D(A).$$

This is precisely the Lagrange identity (3.1). The proof of (3.4) is similar. \square

Remark 3.6. *The identities (3.3) and (3.4) clearly extend to all λ, μ such that $-\lambda, -\bar{\mu} \in \rho(A)$.*

4. THE ŠTRAUS CHARACTERISTIC FUNCTION AND ITS PROPERTIES

The first characteristic function, discussed below, was introduced by Livšić [27]. Later, by completely different methods, a characteristic function was introduced by Sz.-Nagy and Foiaş [44] as part of their harmonic analysis of contractions. As was clarified by M. Krein and Gohberg, the Sz-Nagy-Foiaş characteristic function is a generalisation of the Livšić characteristic function to a wider class of operators. Simultaneously, in a series of papers by Štraus [41, 42], another (unitarily equivalent) characteristic function was introduced in his study of extensions of symmetric operators and also in more general settings. We will introduce the Štraus characteristic function in our setting and discuss its connection to the Sz-Nagy-Foiaş characteristic function below. This definition of the characteristic function goes back to the idea of the characteristic function of an operator knot as introduced by the Odessa school [5]. It is also related to the characteristic functions in the setting of boundary triples, introduced by Derkach and Malamud, see, e.g. [12, 13], where a discussion of different definitions of characteristic functions and their connections can also be found.

We recall that in all of the following A is a maximally dissipative operator on H and Γ, Γ_* and E, E_* are operators and, respectively, spaces with the properties given in Lemma 3.1. We start with a simple identity.

Lemma 4.1. *For all $u \in D(A)$ and $z \in \rho(A^*)$ we have*

$$\|\Gamma_*(A^* - z)^{-1}(A - z)u\|^2 = \|\Gamma u\|^2 - 2\Im(z)\|(A^* - z)^{-1}(A - z)u - u\|^2.$$

Proof. This is an explicit calculation. For $u \in D(A)$ we have, by the second Lagrange identity (3.2),

$$\begin{aligned} &i\langle \Gamma_*(A^* - z)^{-1}(A - z)u, \Gamma_*(A^* - z)^{-1}(A - z)u \rangle_{E_*} \\ &= -\langle A^*(A^* - z)^{-1}(A - z)u, (A^* - z)^{-1}(A - z)u \rangle \\ &\quad + \langle (A^* - z)^{-1}(A - z)u, A^*(A^* - z)^{-1}(A - z)u \rangle \\ &= -\langle (A - z)u, (A^* - z)^{-1}(A - z)u \rangle + \langle (A^* - z)^{-1}(A - z)u, (A - z)u \rangle \\ (4.1) \quad &- 2i\Im(z)\|(A^* - z)^{-1}(A - z)u\|^2 \\ &= -\langle u, (A^* - \bar{z})(A^* - z)^{-1}(A - z)u \rangle + \langle (A^* - \bar{z})(A^* - z)^{-1}(A - z)u, u \rangle \\ &\quad - 2i\Im(z)\|(A^* - z)^{-1}(A - z)u\|^2 \\ &= -\langle u, (A - z)u \rangle + \langle (A - z)u, u \rangle + 2i\Im(z)\langle u, (A^* - z)^{-1}(A - z)u \rangle \\ &\quad + 2i\Im(z)\langle (A^* - z)^{-1}(A - z)u, u \rangle - 2i\Im(z)\|(A^* - z)^{-1}(A - z)u\|^2. \end{aligned}$$

By the first Lagrange identity (3.1), this is equal to

$$i\|\Gamma u\|^2 - 2i\Im(z)\|u\|^2 + 4i\Im(z)\Re\langle u, (A^* - z)^{-1}(A - z)u \rangle - 2i\Im(z)\|(A^* - z)^{-1}(A - z)u\|^2,$$

which, in turn, simplifies to

$$i\left[\|\Gamma u\|^2 - 2\Im(z)\|(A^* - z)^{-1}(A - z)u - u\|^2\right],$$

proving the identity. \square

Corollary 4.2. *For $u \in D(A)$ and $z \in \mathbb{C}^+$,*

$$(4.2) \quad \|\Gamma_*(A^* - z)^{-1}(A - z)u\| \leq \|\Gamma u\|.$$

Hence, there exists a unique contraction $S(z) : E \rightarrow E_$, analytic in the upper half-plane, such that*

$$(4.3) \quad S(z)\Gamma u = \Gamma_*(A^* - z)^{-1}(A - z)u \text{ for all } u \in D(A).$$

Proof. Define $S(z)$ on $\text{Ran}(\Gamma)$ by (4.3). Then $S(z)$ is both well-defined and contractive by (4.2). Therefore, it can be uniquely extended to a contraction on E . Analyticity follows from analyticity of the right hand side of (4.3). \square

Lemma 4.3. For $u \in D(A^*)$ and $z \in \rho(A)$,

$$\|\Gamma(A-z)^{-1}(A^*-z)u\|^2 = \|\Gamma_*u\|^2 + 2\Im(z)\|(A-z)^{-1}(A^*-z)u + u\|^2.$$

Correspondingly, for $z \in \mathbb{C}^-$ there exists a contraction $S_*(z) : E_* \rightarrow E$, analytic in the lower half-plane, such that

$$(4.4) \quad S_*(z)\Gamma_*u = \Gamma(A-z)^{-1}(A^*-z)u.$$

Proof. This is analogous to the proof of Lemma 4.1 and Corollary 4.2. \square

We now wish to extend $S(z)$ by (4.3) to all $z \in \rho(A^*)$ and $S_*(z)$ by (4.4) to all $z \in \rho(A)$.

Lemma 4.4. $S(z)$ is well-defined on $\text{Ran } \Gamma$ for $z \in \rho(A^*)$ and $S_*(z)$ is well-defined on $\text{Ran } \Gamma_*$ for $z \in \rho(A)$.

Proof. We prove the result for $S(z)$, the proof for $S_*(z)$ is similar. We need to show that if $u \in D(A)$ with $\Gamma u = 0$, then for any $z \in \rho(A^*)$ we have $\Gamma_*(A^*-z)^{-1}(A-z)u = 0$. If $\Gamma u = 0$ then by (3.1), we have that $\langle Au, v \rangle = \langle u, Av \rangle$ for any $v \in D(A)$. This implies that $u \in D(A^*)$ with $A^*u = Au$, so $\Gamma_*(A^*-z)^{-1}(A-z)u = \Gamma_*u$. Using (3.2) and again (3.1), we get that

$$-i\|\Gamma_*u\|^2 = \langle A^*u, u \rangle - \langle u, A^*u \rangle = \langle Au, u \rangle - \langle u, Au \rangle = i\|\Gamma u\|^2 = 0,$$

as required. \square

The following lemma gives a useful identity for the difference of S at two different points.

Lemma 4.5. For $\mu, \tilde{\mu} \in \rho(A^*)$, we have the following identity:

$$(4.5) \quad S(\mu) - S(\tilde{\mu}) = i(\mu - \tilde{\mu}) (\Gamma_*(A^* - \mu)^{-1}) \left(\Gamma(A - \tilde{\mu})^{-1} \right)^* \text{ on } \text{Ran } \Gamma.$$

Proof. Let $h \in D(A)$ and set

$$v = \left[S(\mu) - S(\tilde{\mu}) - i(\mu - \tilde{\mu}) (\Gamma_*(A^* - \mu)^{-1}) \left(\Gamma(A - \tilde{\mu})^{-1} \right)^* \right] \Gamma h.$$

Then using (4.3) we get

$$\begin{aligned} v &= \Gamma_*(A^* - \mu)^{-1}(A - \mu)h - \Gamma_*(A^* - \tilde{\mu})^{-1}(A - \tilde{\mu})h + i(\tilde{\mu} - \mu) (\Gamma_*(A^* - \mu)^{-1}) \left(\Gamma(A - \tilde{\mu})^{-1} \right)^* \Gamma h \\ &= \Gamma_*(A^* - \mu)^{-1} \left[(A - \mu)h - (A - \tilde{\mu})h - (\tilde{\mu} - \mu)(A^* - \tilde{\mu})^{-1}(A - \tilde{\mu})h + i(\tilde{\mu} - \mu) \left(\Gamma(A - \tilde{\mu})^{-1} \right)^* \Gamma h \right] \\ &= (\tilde{\mu} - \mu)\Gamma_*(A^* - \mu)^{-1} \left[h - (A^* - \tilde{\mu})^{-1}(A - \tilde{\mu})h + i \left(\Gamma(A - \tilde{\mu})^{-1} \right)^* \Gamma h \right] \end{aligned}$$

Next, set $h = (A^* - \tilde{\mu})^{-1}g$. After applying (3.3) (with Remark 3.6), a short calculation shows that the term in the square brackets vanishes, giving $v = 0$, as required. \square

From this identity we see that although $S(\mu)$ need not be a contraction for all $\mu \in \rho(A^*)$, it remains bounded on $\text{Ran } \Gamma$.

Corollary 4.6. For $\mu \in \rho(A^*)$, we have that $S(\mu)$ is a bounded operator on $\text{Ran } \Gamma$. Moreover,

$$(4.6) \quad \|S(\mu)\| \leq 1 + 16\sqrt{\gamma(1+\gamma)},$$

where $\gamma = |\Im(\mu)| \|(A^* - \mu)^{-1}\|$.

Proof. Choose $\tilde{\mu} \in \mathbb{C}^+$. Then from (4.5), we get that

$$\begin{aligned} \|S(\mu)\| &\leq \|S(\tilde{\mu})\| + |\mu - \tilde{\mu}| \|\Gamma_*(A^* - \mu)^{-1}\| \|\Gamma(A - \tilde{\mu})^{-1}\| \\ &\leq 1 + |\mu - \tilde{\mu}| \|\Gamma_*(A^* - \mu)^{-1}\| \|\Gamma(A - \tilde{\mu})^{-1}\|. \end{aligned}$$

Let $A' = A - \Re(\tilde{\mu})$. Then, using the Lagrange identity (3.1),

$$\begin{aligned} \|\Gamma(A - \tilde{\mu})^{-1}\|^2 &\leq 2 \|(A - \tilde{\mu})^{-1}\| \|A'(A - \tilde{\mu})^{-1}\| \\ &\leq \frac{2}{\Im(\tilde{\mu})} \|I + i\Im(\tilde{\mu})(A - \tilde{\mu})^{-1}\| \leq \frac{4}{\Im(\tilde{\mu})}. \end{aligned}$$

Next, let $\lambda = \Re(\mu) + i\tau$ form some $\tau > 0$. Then using the previous estimate, we get

$$\begin{aligned} \|\Gamma_*(A^* - \mu)^{-1}\| &\leq \|\Gamma_*(A^* - \lambda)^{-1} [I + (\mu - \lambda)(A^* - \mu)^{-1}]\| \\ &\leq \|\Gamma_*(A^* - \lambda)^{-1}\| \|I + i(\Im(\mu) - \tau)(A^* - \mu)^{-1}\| \\ &\leq \sqrt{\frac{4}{\tau}} [1 + (|\Im(\mu)| + \tau) \|(A^* - \mu)^{-1}\|]. \end{aligned}$$

Combining the estimates, we get that

$$\|S(\mu)\| \leq 1 + \frac{4(|\Im(\mu)| + \Im(\tilde{\mu}))}{\sqrt{\Im\tilde{\mu}}} \cdot \frac{1 + (|\Im(\mu)| + \tau) \|(A^* - \mu)^{-1}\|}{\sqrt{\tau}}.$$

Both fractions are of the form $(a + bx)/\sqrt{x}$ and are minimized for $x = a/b$ with value $2\sqrt{ab}$. Thus,

$$\|S(\mu)\| \leq 1 + 16\sqrt{\Im\mu} \cdot \sqrt{(1 + |\Im(\mu)|) \|(A^* - \mu)^{-1}\|} \|(A^* - \mu)^{-1}\|,$$

as claimed. \square

This justifies the following definition.

Definition 4.7. *The operator-valued function $S(\cdot)$, defined for $z \in \rho(A^*)$ by (4.3) on $\text{Ran}(\Gamma)$ and extended to E by continuity is called the Štraus characteristic function of the operator A .*

Lemma 4.8. $S(z) = S_*(\bar{z})$ for $z \in \rho(A^*)$.

Proof. We need to show that for all $e \in E$ and $e_* \in E_*$, we have

$$\langle S(z)e, e_* \rangle_{E_*} = \langle e, S_*(\bar{z})e_* \rangle_E.$$

Due to the boundedness of $S(z)$ and $S_*(\bar{z})$, it is sufficient to show this on dense sets. Therefore, we choose $e = \Gamma(A - \bar{z})^{-1}f$ and $e_* = \Gamma_*(A^* - z)^{-1}g$ for $f, g \in H$. Then

$$i \langle S(z)e, e_* \rangle_{E_*} = i \langle \Gamma_*(A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}f, \Gamma_*(A^* - z)^{-1}g \rangle_{E_*}$$

and using the Lagrange identity (3.2) this equals

$$\begin{aligned} & - \langle (A^* - \bar{z})(A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}f, (A^* - z)^{-1}g \rangle + \langle (A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}f, g \rangle \\ & = \langle (I - (A - \bar{z})^{-1}(A^* - \bar{z})) (A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}f, g \rangle \\ & = \langle ((A^* - z)^{-1} - (A - \bar{z})^{-1} (I + (z - \bar{z})(A^* - z)^{-1})) (A - z)(A - \bar{z})^{-1}f, g \rangle \\ & = \langle ((A^* - z)^{-1} - (A - \bar{z})^{-1} - (z - \bar{z})(A - \bar{z})^{-1}(A^* - z)^{-1}) (A - z)(A - \bar{z})^{-1}f, g \rangle. \end{aligned}$$

On the other hand, using the first Lagrange identity (3.1) we have

$$\begin{aligned} i \langle e, S_*(\bar{z})e_* \rangle_E &= i \langle \Gamma(A - \bar{z})^{-1}f, \Gamma(A - \bar{z})^{-1}(A^* - \bar{z})(A^* - z)^{-1}g \rangle \\ &= \langle (A - z)(A - \bar{z})^{-1}f, (A - \bar{z})^{-1}(A^* - \bar{z})(A^* - z)^{-1}g \rangle - \langle (A - \bar{z})^{-1}f, (A^* - \bar{z})(A^* - z)^{-1}g \rangle \\ &= \langle (A - z)(A - \bar{z})^{-1}f, ((A - \bar{z})^{-1}(A^* - \bar{z}) - I)(A^* - z)^{-1}g \rangle \\ &= \langle (A - z)(A - \bar{z})^{-1}f, ((A - \bar{z})^{-1}(I + (z - \bar{z})(A^* - z)^{-1}) - (A^* - z)^{-1})g \rangle \\ &= \langle (A - z)(A - \bar{z})^{-1}f, ((A - \bar{z})^{-1} - (A^* - z)^{-1} + (z - \bar{z})(A - \bar{z})^{-1}(A^* - z)^{-1})g \rangle \\ &= \langle ((A^* - z)^{-1} - (A - \bar{z})^{-1} - (z - \bar{z})(A - \bar{z})^{-1}(A^* - z)^{-1}) (A - z)(A - \bar{z})^{-1}f, g \rangle, \end{aligned}$$

proving the desired equality. \square

Lemma 4.9. $S(z)S_*(z) = I_{E_*}$ and $S_*(z)S(z) = I_E$ whenever $z \in \rho(A) \cap \rho(A^*)$.

Proof. Due to the boundedness of the operators involved, it is again sufficient to show this on a dense set. Let $u \in D(A^*)$. Then

$$\begin{aligned} S(z)S_*(z)\Gamma_*u &= S(z)\Gamma(A - z)^{-1}(A^* - z)u \\ &= \Gamma_*(A^* - z)^{-1}(A - z)(A - z)^{-1}(A^* - z)u = \Gamma_*u. \end{aligned}$$

The second equality can be proved similarly. \square

This immediately gives the following results.

Corollary 4.10. (1) $S(z)$ is unitary for $z \in \mathbb{R} \cap \rho(A)$.

(2) If $\sigma(A)$ does not cover the whole upper half plane (or if $\rho(A) \cap \rho(A^*) \neq \emptyset$), then $\dim E = \dim E_*$.

Proof. (1) This is immediate from Lemmas 4.8 and 4.9.

(2) Let $z \in \rho(A) \cap \rho(A^*)$. From the first equality in Lemma 4.9, we get that

$$\dim E_* \leq \text{rank } S(z) \leq \min\{\dim E, \dim E_*\},$$

while the second equality gives

$$\dim E \leq \text{rank } S_*(z) \leq \min\{\dim E, \dim E_*\},$$

proving the result. \square

The next lemma shows Hermitian positivity properties (see Azizov-Iokhvidov [1] for related results).

Lemma 4.11. *For $w, z \in \mathbb{C}^+$, we have*

$$\frac{1}{\bar{w} - z} (I_E - S^*(w)S(z)) = i (\Gamma(A - \bar{w})^{-1}) (\Gamma(A - \bar{z})^{-1})^*$$

and for $w, z \in \mathbb{C}^-$, we have

$$\frac{1}{\bar{w} - z} (I_{E_*} - S_*^*(w)S_*(z)) = -i (\Gamma_*(A^* - \bar{w})^{-1}) (\Gamma_*(A^* - \bar{z})^{-1})^*.$$

Proof. We check the first equality on a dense set. Let $u = \Gamma(A - \bar{z})^{-1}f$ for some $f \in H$. Then using that $S^*(w) = S_*(\bar{w})$ from Lemma 4.8, we get

$$\begin{aligned} (I_E - S^*(w)S(z))u &= [\Gamma(A - \bar{z})^{-1} - S^*(w)\Gamma_*(A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}]f \\ &= [\Gamma(A - \bar{z})^{-1} - \Gamma(A - \bar{w})^{-1}(A^* - \bar{w})(A^* - z)^{-1}(A - z)(A - \bar{z})^{-1}]f \\ &= \Gamma(A - \bar{w})^{-1} [I + (\bar{z} - \bar{w})(A - \bar{z})^{-1} - (I + (z - \bar{w})(A^* - z)^{-1})(I + (\bar{z} - z)(A - \bar{z})^{-1})]f \\ &= (z - \bar{w})\Gamma(A - \bar{w})^{-1} [(A - \bar{z})^{-1} - (A^* - z)^{-1} - (A^* - z)^{-1}(\bar{z} - z)(A - \bar{z})^{-1}]f. \end{aligned}$$

Using (3.3) with $\lambda = -\bar{z}$ and $\mu = -z$, we get

$$\begin{aligned} (I_E - S^*(w)S(z))u &= (z - \bar{w})\Gamma(A - \bar{w})^{-1}(-i)(\Gamma(A - \bar{z})^{-1})^*(\Gamma(A - \bar{z})^{-1})f \\ &= i(\bar{w} - z) (\Gamma(A - \bar{w})^{-1}) (\Gamma(A - \bar{z})^{-1})^* u, \end{aligned}$$

proving the identity. \square

Remark 4.12. *In the case when $z = w$, the rank of the limit operator of $I_E - S^*(z)S(z)$ as z tends to the real axis corresponds to the local multiplicity of the a.c.-spectrum of A (see [35, 45]).*

For later calculations, we will also need the following identities:

Lemma 4.13. *For any $u \in H$, $\mu, z \in \mathbb{C}^-$ we have*

$$(4.7) \quad (\Gamma_*(A^* - \bar{\mu})^{-1})^* S(\bar{z}) = [I - (\bar{z} - \mu)(A - \mu)^{-1}] (\Gamma(A - z)^{-1})^*$$

and

$$(4.8) \quad (\Gamma(A - \mu)^{-1})^* S_*(z) = [I - (z - \bar{\mu})(A^* - \bar{\mu})^{-1}] (\Gamma_*(A^* - \bar{z})^{-1})^*.$$

Proof. From the definition of S_* in (4.4) we have that

$$S_*(z)\Gamma_*(A^* - \bar{\mu})^{-1} = \Gamma(A - z)^{-1}(A^* - z)(A^* - \bar{\mu})^{-1}.$$

Taking adjoints, using Lemma 4.8, we get

$$\begin{aligned} (\Gamma_*(A^* - \bar{\mu})^{-1})^* S(\bar{z}) &= [(A^* - z)(A^* - \bar{\mu})^{-1}]^* (\Gamma(A - z)^{-1})^* \\ &= [I - (z - \bar{\mu})(A^* - \bar{\mu})^{-1}]^* (\Gamma(A - z)^{-1})^* \\ &= [I - (\bar{z} - \mu)(A - \mu)^{-1}] (\Gamma(A - z)^{-1})^*. \end{aligned}$$

This proves (4.7). The proof of (4.8) is similar. \square

Example 4.14. *As in Example 3.4, we consider the Schrödinger operator on the half-line \mathbb{R}^+ with a dissipative boundary condition and potential. We have seen that for $u, v \in D(A)$, equation (3.1) holds with*

$$\Gamma u = \begin{pmatrix} \sqrt{2\Im h} u(0) \\ \sqrt{2\Im q} u \end{pmatrix},$$

which has dense range in $E = E_* = \mathbb{C} \oplus L^2(\{x \in \mathbb{R}^+ : \Im q(x) > 0\})$. Note that A^* is given by

$$(A^* f)(x) = -f''(x) + \overline{q(x)}f(x) \quad \text{with domain} \quad D(A^*) := \{y \in H^2(\mathbb{R}^+) : y'(0) = \bar{h}y(0)\}.$$

Γ_* is given by the same expression as Γ .

We now calculate the characteristic function. For $u \in D(A)$ consider $S(z)\Gamma u = \Gamma_*(A^* - z)^{-1}(A - z)u$ with $z \in \rho(A^*)$. We choose φ_* and ψ_* to be the fundamental solutions of $-y'' + \bar{q}y = \lambda y$ with φ_* and ψ_* normalized by

$$(4.9) \quad \begin{pmatrix} \varphi_*(0) \\ \varphi'_*(0) \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} \psi_*(0) \\ \psi'_*(0) \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}.$$

Moreover, let m_* denote the Weyl-Titchmarsh function associated with $-y'' + \bar{q}y$, i.e. $m_*(z)\varphi_* + \psi_*$ is the L^2 -solution to $-y'' + \bar{q}y = zy$. This solution is unique up to constants due to our assumptions on q . Choosing

$$(4.10) \quad c_*(z) = \frac{h - \bar{h}}{\bar{h} - m_*(z)} u(0),$$

we have that $v = u + c_*(z)(m_*(z)\varphi_* + \psi_*) + 2i(A^* - z)^{-1}(\Im q) u$, $v \in D(A^*)$ and for any $u \in D(A)$,

$$(A^* - z)v = -u'' + (\bar{q} - z)u + 2i(\Im q) u = (A - z)u.$$

Hence, $v = (A^* - z)^{-1}(A - z)u$ and for $z \in \rho(A^*)$ we have $S(z)\Gamma u = \Gamma_* v$, i.e.

$$S(z) \begin{pmatrix} \sqrt{2\Im h} u(0) \\ \sqrt{2\Im q} u \end{pmatrix} = \begin{pmatrix} \sqrt{2\Im h} v(0) \\ \sqrt{2\Im q} v \end{pmatrix} = \begin{pmatrix} \sqrt{2\Im h} (u(0) + c_*(z) + (2i(A^* - z)^{-1}(\Im q) u)(0)) \\ \sqrt{2\Im q} (u + c_*(z)(m_*(z)\varphi_* + \psi_*) + 2i(A^* - z)^{-1}(\Im q) u) \end{pmatrix}.$$

This implies that

$$(4.11) \quad S(z) = \begin{pmatrix} I + \frac{h - \bar{h}}{h - m_*(z)} & i\sqrt{2\Im h} (A^* - z)^{-1} \sqrt{2\Im q} \cdot (0) \\ i \frac{\sqrt{2\Im h} \sqrt{2\Im q}}{h - m_*(z)} (m_*(z)\varphi_*(x) + \psi_*(x)) & I + i\sqrt{2\Im q} (A^* - z)^{-1} \sqrt{2\Im q} \end{pmatrix} \\ = \begin{pmatrix} \frac{h - m_*(z)}{h - m_*(z)} & i\sqrt{2\Im h} \int_0^\infty \frac{m_*(z)\varphi_*(y) + \psi_*(y)}{h - m_*(z)} \sqrt{2\Im q(y)} \cdot (y) dy \\ i \frac{\sqrt{2\Im h} \sqrt{2\Im q}}{h - m_*(z)} (m_*(z)\varphi_*(x) + \psi_*(x)) & I + i\sqrt{2\Im q} (A^* - z)^{-1} \sqrt{2\Im q} \end{pmatrix}.$$

Here, the top right entry is a map from $L^2(\Omega)$ to \mathbb{C} and $\cdot(y)$ indicates that the L^2 -function needs to be inserted for \cdot and then evaluated.

Note that the top left entry is the same as the well-known formula of Pavlov for the case of real q , where the functions m and m_* coincide:

$$(4.12) \quad S(z) = \left(1 + \frac{h - \bar{h}}{\bar{h} - m(z)} \right).$$

Pavlov deduced it from the scattering theory interpretation of the characteristic function [34, 36], for a different approach see also [12]. The bottom right entry agrees with the Livšič characteristic function for the case with a selfadjoint boundary condition [27]. Moreover, this formula shows the connection between the Weyl m -function and the characteristic function for this example.

We conclude this section by showing that the Štraus characteristic function as defined here coincides up to an isometric transformation with the Sz-Nagy-Foias characteristic function for a contraction T , see [44], given by

$$(4.13) \quad \Theta(\lambda) = [-T + \lambda D_{T^*} (I - \lambda T^*)^{-1} D_T] |_{\text{Ran } \overline{D_T}}, \quad \text{for } |\lambda| < 1,$$

where $D_T = \sqrt{I - T^*T}$ and $D_{T^*} = \sqrt{I - TT^*}$. Then $\Theta(\lambda) \in B(\text{Ran}(D_T), \overline{\text{Ran}(D_{T^*})})$.

Proposition 4.15. *Let $z \in \mathbb{C}^+$ and A be a maximal dissipative operator with Cayley transform T given by (2.1). Then there exist isometric surjective operators $U : E \rightarrow \text{Ran}(D_T)$ and $U_* : E_* \rightarrow \text{Ran}(D_{T^*})$ such that*

$$\Theta(\lambda) = U_* S(z) U^*, \quad \text{where } \lambda = \frac{z - i}{z + i}.$$

Proof. We first determine some of the expressions arising in $\Theta(\lambda)$ in terms of A and A^* . We have

$$(4.14) \quad \begin{aligned} D_T^2 &= I - T^*T = I - (I + 2i(A^* - i)^{-1})(I - 2i(A + i)^{-1}) \\ &= 2i[(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1}(A + i)^{-1}]. \end{aligned}$$

Similarly,

$$(4.15) \quad D_{T^*}^2 = 2i[(A + i)^{-1} - (A^* - i)^{-1} + 2i(A + i)^{-1}(A^* - i)^{-1}].$$

Moreover,

$$\begin{aligned} I - \lambda T^* &= I - \lambda - 2i\lambda(A^* - i)^{-1} = ((1 - \lambda)(A^* - i) - 2i\lambda)(A^* - i)^{-1} \\ &= (1 - \lambda) \left(A^* - i - \frac{2i\lambda}{1 - \lambda} \right) (A^* - i)^{-1} = (1 - \lambda)(A^* - z)(A^* - i)^{-1}, \end{aligned}$$

so

$$(4.16) \quad (I - \lambda T^*)^{-1} = \frac{1}{1 - \lambda} (A^* - i)(A^* - z)^{-1}.$$

From (3.1), we have that for any $f, g \in H$,

$$\begin{aligned} i \langle \Gamma(A + i)^{-1} f, \Gamma(A + i)^{-1} g \rangle_E &= \langle A(A + i)^{-1} f, (A + i)^{-1} g \rangle_H - \langle (A + i)^{-1} f, A(A + i)^{-1} g \rangle_H \\ &= \langle [(A^* - i)^{-1} A(A + i)^{-1} - A^*(A^* - i)^{-1} (A + i)^{-1}] f, g \rangle_H \\ &= \langle [(A^* - i)^{-1} - (A + i)^{-1} - 2i(A^* - i)^{-1} (A + i)^{-1}] f, g \rangle_H \\ &= \left\langle \frac{i}{2} D_T^2 f, g \right\rangle_H. \end{aligned}$$

This implies that $\frac{1}{2} D_T^2 = (\Gamma(A + i)^{-1})^* \Gamma(A + i)^{-1}$. Similarly, $\frac{1}{2} D_{T^*}^2 = (\Gamma_*(A^* - i)^{-1})^* \Gamma_*(A^* - i)^{-1}$. Therefore there exist U and U_* with the desired properties, such that

$$D_T = \sqrt{2} U \Gamma(A + i)^{-1} \quad \text{and} \quad D_{T^*} = -\sqrt{2} U_* \Gamma_*(A^* - i)^{-1}.$$

Noting that $T D_T^2 = D_{T^*}^2 T$ and thus, using the functional calculus $T D_T = D_{T^*} T$, we now have with (4.14) and (4.16) that

$$\begin{aligned} \Theta(\lambda) D_T &= -D_{T^*} T + \frac{2i\lambda}{1 - \lambda} D_{T^*} (A^* - i)(A^* - z)^{-1} [(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1} (A + i)^{-1}] \\ &= D_{T^*} \left(-I + 2i(A + i)^{-1} + (z - i)(A^* - i)(A^* - z)^{-1} [(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1} (A + i)^{-1}] \right) \\ &= D_{T^*} (A^* - i)(A^* - z)^{-1} \\ &\quad \left[(I - (z - i)(A^* - i)^{-1})(-I + 2i(A + i)^{-1}) + (z - i) [(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1} (A + i)^{-1}] \right] \\ &= D_{T^*} (A^* - i)(A^* - z)^{-1} (-I + (z + i)(A + i)^{-1}). \end{aligned}$$

Using our expressions for D_{T^*} and D_T , we get

$$\begin{aligned} \Theta(\lambda) D_T &= \sqrt{2} U_* \Gamma_*(A^* - z)^{-1} (-I + (z + i)(A + i)^{-1}) \\ &= \sqrt{2} U_* \Gamma_*(A^* - z)^{-1} (A - z)(A + i)^{-1} \\ &= \sqrt{2} U_* S(z) \Gamma(A + i)^{-1} = U_* S(z) U^* D_T, \end{aligned}$$

proving that $\Theta(\lambda) = U_* S(z) U^*$. \square

Remark 4.16. *This shows that the two approaches differ in the choice of the root of the imaginary part of the operator. The advantage of the Štraus characteristic function for us is that we can often explicitly determine Γ and Γ_* , while it is rarely possible to find explicit expressions for D_T and D_{T^*} .*

5. DEFINITION OF THE DILATION

Before studying the dilation itself, we first introduce its domain and show that it has several equivalent descriptions. Here and in what follows we let $\mathcal{H} = L^2(\mathbb{R}^-, E_*) \oplus H \oplus L^2(\mathbb{R}^+, E)$, where $L^2(\mathbb{R}^-, E_*)$ and $L^2(\mathbb{R}^+, E)$ are suitable channels in the sense of Lax and Phillips [24].

Definition 5.1. *Let $\mu \in \mathbb{C}^-$ and $\lambda \in \mathbb{C}^+$. We define the subset $D(\mathcal{L})$ of \mathcal{H} by*

$$(5.1) \quad D(\mathcal{L}) = \left\{ U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} : u \in H, v_+ \in H^1(\mathbb{R}^+, E), v_- \in H^1(\mathbb{R}^-, E_*), \right. \\ \left. \begin{aligned} u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) &\in D(A), u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \in D(A^*), \\ (I) v_+(0) &= S^*(-\mu) v_-(0) + i\Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)), \\ (II) v_-(0) &= S(-\bar{\lambda}) v_+(0) - i\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^* v_+(0)) \end{aligned} \right\}.$$

We will see in Lemma 5.4 that the conditions (I) and (II) are equivalent, so it is possible to omit one of them in the definition. The numbers μ and λ are regularisation parameters. For special situations, such as if the imaginary part of A is relatively bounded, they are not needed. For general MDOs, however, the regularisation is necessary. Besides the disadvantage of complicating the expressions in the boundary conditions (I) and (II), the presence of

the parameters also may bring some advantages, e.g. allowing to simplify the conditions by a particular choice of the parameters or by taking limits whenever the terms in (5.1) admit suitable asymptotics.

We now show that $D(\mathcal{L})$ is independent of the choice of $\mu \in \mathbb{C}^-$ and $\lambda \in \mathbb{C}^+$. First, we show this for the conditions to lie in the domains of the operators, and in a second lemma we consider (I) and (II).

Lemma 5.2. *Let $u \in H$, $v_-(0) \in E_*$ and $v_+(0) \in E$. If $u + (\Gamma_*(A^* + \mu_0)^{-1})^* v_-(0) \in D(A)$ for some $\mu_0 \in \mathbb{C}^-$, then $u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \in D(A)$ for all $\mu \in \mathbb{C}^-$. Similarly, if $u + (\Gamma(A + \lambda_0)^{-1})^* v_+(0) \in D(A^*)$ for some $\lambda_0 \in \mathbb{C}^+$, then $u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \in D(A^*)$ for all $\lambda \in \mathbb{C}^+$.*

Proof. We show that $u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \in D(A)$ for all $\mu \in \mathbb{C}^-$. Clearly, it suffices to show that

$$\left((\Gamma_*(A^* + \mu)^{-1})^* - (\Gamma_*(A^* + \mu_0)^{-1})^* \right) v_-(0) = \left[\Gamma_*((A^* + \mu)^{-1} - (A^* + \mu_0)^{-1}) \right]^* v_-(0) \in D(A).$$

We have

$$\begin{aligned} \left[\Gamma_*((A^* + \mu)^{-1} - (A^* + \mu_0)^{-1}) \right]^* v_-(0) &= [(\mu_0 - \mu) \Gamma_*((A^* + \mu)^{-1} (A^* + \mu_0)^{-1})]^* v_-(0) \\ &= (\bar{\mu}_0 - \bar{\mu}) (A + \bar{\mu}_0)^{-1} (\Gamma_*(A^* + \mu)^{-1})^* v_-(0), \end{aligned}$$

which clearly lies in $D(A)$. □

Now we are in a position to prove independence of the boundary conditions (I) and (II) from the choice of $\mu \in \mathbb{C}^-$ and $\lambda \in \mathbb{C}^+$.

Lemma 5.3. *The conditions (I) and (II) in (5.1) are independent of $\mu \in \mathbb{C}^-$ and $\lambda \in \mathbb{C}^+$, respectively.*

Proof. We check this for (I). We need to show that for any $\mu, \tilde{\mu} \in \mathbb{C}^-$ we have

$$(5.2) \quad S^*(-\mu) - S^*(-\tilde{\mu}) + i\Gamma \left(\Gamma_*((A^* + \mu)^{-1} - (A^* + \tilde{\mu})^{-1}) \right)^* = 0 \text{ on } E_*.$$

Since,

$$\begin{aligned} \Gamma \left(\Gamma_*((A^* + \mu)^{-1} - (A^* + \tilde{\mu})^{-1}) \right)^* &= \Gamma \left((\tilde{\mu} - \mu) \Gamma_*((A^* + \mu)^{-1} (A^* + \tilde{\mu})^{-1}) \right)^* \\ &= (\bar{\tilde{\mu}} - \bar{\mu}) \left(\Gamma(A + \bar{\tilde{\mu}})^{-1} \right) (\Gamma_*(A^* + \mu)^{-1})^*, \end{aligned}$$

the result follows from taking adjoints in (4.5). □

The next lemma shows that the conditions (I) and (II) in (5.1) are equivalent.

Lemma 5.4. *Let $\mu \in \mathbb{C}^-$, $\lambda \in \mathbb{C}^+$, $u \in H$, $v_+(0) \in E$ and $v_-(0) \in E_*$. Then $u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \in D(A)$ and*

$$(I) \quad v_+(0) = S^*(-\mu) v_-(0) + i\Gamma \left(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \right)$$

holds if and only if $u + (\Gamma(A + \lambda)^{-1})^ v_+(0) \in D(A^*)$ and*

$$(II) \quad v_-(0) = S(-\bar{\lambda}) v_+(0) - i\Gamma_* \left(u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \right).$$

Proof. We will assume that $u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \in D(A^*)$ and (II) holds. The proof of the converse is similar. First, we need to show that $w := u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \in D(A)$. Now,

$$w = u + (\Gamma_*(A^* + \mu)^{-1})^* \left[S(-\bar{\lambda}) v_+(0) - i\Gamma_* \left(u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \right) \right].$$

Using (4.7), we get

$$w = u + \left[I + (\bar{\lambda} - \bar{\mu})(A + \bar{\mu})^{-1} \right] (\Gamma(A + \lambda)^{-1})^* v_+(0) - i(\Gamma_*(A^* + \mu)^{-1})^* \Gamma_* \left(u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \right).$$

Clearly, $w \in D(A)$ if and only if $\tilde{w} \in D(A)$, where

$$\begin{aligned} \tilde{w} &:= u + (\Gamma(A + \lambda)^{-1})^* v_+(0) - i(\Gamma_*(A^* + \mu)^{-1})^* \Gamma_* \left(u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \right) \\ &= \left[I - i(\Gamma_*(A^* + \mu)^{-1})^* \Gamma_* \right] \left(u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \right). \end{aligned}$$

Next, using that $u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \in D(A^*)$, for $\nu \in \mathbb{C}^+$ there exists f such that

$$u + (\Gamma(A + \lambda)^{-1})^* v_+(0) = (A^* - \nu)^{-1} f.$$

Inserting this in \tilde{w} and using (3.4), we get

$$\begin{aligned} \tilde{w} &= (A^* - \nu)^{-1} f + \left[(A + \bar{\mu})^{-1} - (A^* - \nu)^{-1} + (\bar{\mu} + \nu)(A + \bar{\mu})^{-1} (A^* - \nu)^{-1} \right] f \\ &= \left[(A + \bar{\mu})^{-1} + (\bar{\mu} + \nu)(A + \bar{\mu})^{-1} (A^* - \nu)^{-1} \right] f \in D(A). \end{aligned}$$

It remains to check (I), which using Lemma 4.8, is equivalent to

$$(5.3) \quad v_+(0) - S^*(-\mu)S(-\bar{\lambda})v_+(0) = -iS_*(-\bar{\mu})\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)) \\ + i\Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^*(S(-\bar{\lambda})v_+(0) - i\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^*v_+(0))))).$$

By Lemma 4.11, the left hand side of (5.3) is given by

$$(5.4) \quad i(\bar{\lambda} - \bar{\mu})(\Gamma(A + \bar{\mu})^{-1})(\Gamma(A + \lambda)^{-1})^*v_+(0).$$

We now calculate the terms on the right hand side of (5.3). By the definition of S_* , we have

$$(5.5) \quad S_*(-\bar{\mu})\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)) = \Gamma(A + \bar{\mu})^{-1}(A^* + \bar{\mu})(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)),$$

by (4.7) we have

$$(5.6) \quad (\Gamma_*(A^* + \mu)^{-1})^*S(-\bar{\lambda})v_+(0) = [I + (\bar{\lambda} - \bar{\mu})(A + \bar{\mu})^{-1}](\Gamma(A + \lambda)^{-1})^*v_+(0) \\ = (\Gamma(A + \lambda)^{-1})^*v_+(0) + (\bar{\lambda} - \bar{\mu})(A + \bar{\mu})^{-1}(\Gamma(A + \lambda)^{-1})^*v_+(0).$$

Applying $i\Gamma$ to the second term (which lies in $D(A)$) gives the same as in (5.4), so this will precisely cancel the left hand side in (5.3). The remaining terms on the right hand side of (5.3) now equal

$$-i\Gamma(A + \bar{\mu})^{-1}(A^* + \bar{\mu})(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)) \\ + i\Gamma(u + (\Gamma(A + \lambda)^{-1})^*v_+(0) - i(\Gamma_*(A^* + \mu)^{-1})^*\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^*v_+(0))).$$

Setting $u + (\Gamma(A + \lambda)^{-1})^*v_+(0) = (A^* + \bar{\lambda})^{-1}f$, this is

$$-i\Gamma[(A + \bar{\mu})^{-1}(A^* + \bar{\mu})(A^* + \bar{\lambda})^{-1}f - (A^* + \bar{\lambda})^{-1}f + i(\Gamma_*(A^* + \mu)^{-1})^*(\Gamma_*(A^* + \bar{\lambda})^{-1}f)].$$

Applying (3.4) shows that this equals zero, as required. \square

Remark 5.5. *This shows which are the free parameters in the description of $D(\mathcal{L})$:*

- *The value for $v_+(0) \in E$,*
- *a vector $h \in D(A^*)$ (such that $u = h - (\Gamma(A + \lambda)^{-1})^*v_+(0)$) and*
- *two vector-valued functions $w_{\pm} \in H^1(\mathbb{R}^{\pm})$ with $w_+(0) = 0 = w_-(0)$.*

So far, we have discussed the domain of the dilation. We now present two equivalent formulae for the part of the action of the dilation on the original Hilbert space H .

Definition 5.6. *Let $\mu \in \mathbb{C}^-$ and $\lambda \in \mathbb{C}^+$. For $U \in D(\mathcal{L})$ as given in (5.1), define two operators $T, T_* : D(\mathcal{L}) \rightarrow H$ by*

$$(5.7) \quad TU := A^*(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)) + \bar{\lambda}(\Gamma(A + \lambda)^{-1})^*v_+(0)$$

and

$$(5.8) \quad T_*U := A(u + (\Gamma_*(A^* + \mu)^{-1})^*v_-(0)) + \bar{\mu}(\Gamma_*(A^* + \mu)^{-1})^*v_-(0).$$

Note that we have

$$(5.9) \quad TU = (A^* + \bar{\lambda})(u + (\Gamma(A + \lambda)^{-1})^*v_+(0)) - \bar{\lambda}u$$

and

$$(5.10) \quad T_*U = (A + \bar{\mu})(u + (\Gamma_*(A^* + \mu)^{-1})^*v_-(0)) - \bar{\mu}u.$$

Lemma 5.7. *The operators T and T_* coincide on $D(\mathcal{L})$.*

Proof. Let $U \in D(\mathcal{L})$, $\lambda \in \mathbb{C}^+$ and $\mu \in \mathbb{C}^-$. We test the equality with functions g from the dense set $D(A)$:

$$\begin{aligned} \langle g, TU \rangle &= \langle Ag, u \rangle + \langle \Gamma(A + \lambda)^{-1}Ag, v_+(0) \rangle + \langle \lambda\Gamma(A + \lambda)^{-1}g, v_+(0) \rangle \\ &= \langle Ag, u \rangle + \langle \Gamma g, v_+(0) \rangle \\ &= \langle Ag, u + (\Gamma_*(A^* + \mu)^{-1})^*v_-(0) \rangle - \langle Ag, (\Gamma_*(A^* + \mu)^{-1})^*v_-(0) \rangle + \langle \Gamma g, v_+(0) \rangle. \end{aligned}$$

Using the Lagrange identity (3.1) for the first term gives

$$\begin{aligned}
\langle g, TU \rangle &= \langle g, A(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) \rangle + i \langle \Gamma g, \Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) \rangle \\
&\quad - \langle Ag, (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \rangle + \langle \Gamma g, v_+(0) \rangle \\
&= \langle g, T_* U - \bar{\mu} (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \rangle + i \langle \Gamma g, \Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) \rangle \\
&\quad - \langle Ag, (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \rangle + \langle \Gamma g, v_+(0) \rangle \\
&= \langle g, T_* U \rangle - \mu \langle g, (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \rangle + \langle \Gamma g, v_+(0) - i \Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) \rangle \\
&\quad - \langle Ag, (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \rangle \\
&= \langle g, T_* U \rangle - \langle (\Gamma_*(A^* + \mu)^{-1})(A + \mu)g, v_-(0) \rangle + \langle \Gamma g, S^*(-\mu)v_-(0) \rangle,
\end{aligned}$$

where the last equality follows from the boundary condition (I) in $D(\mathcal{L})$. Therefore, we get

$$\langle g, TU \rangle - \langle g, T_* U \rangle = \langle S(-\mu)\Gamma g - (\Gamma_*(A^* + \mu)^{-1})(A + \mu)g, v_-(0) \rangle = 0,$$

by definition of the characteristic function. Hence $TU = T_* U$. \square

The previous result immediately shows the following corollary, which justifies the absence of the parameters λ and μ in our notation of T and T_* .

Corollary 5.8. *T and T_* on $D(\mathcal{L})$ are independent of λ and μ .*

Finally, having defined the domain on the dilation $D(\mathcal{L})$ in Definition 5.1, we can now give its full action.

Definition 5.9. *We define the operator \mathcal{L} on \mathcal{H} with domain $D(\mathcal{L})$ by*

$$\mathcal{L}U = \mathcal{L} \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = \begin{pmatrix} iv'_- \\ TU \\ iv'_+ \end{pmatrix}.$$

We see that in the so-called incoming and outgoing channels (the first and last components), see, e.g. [24], the operator \mathcal{L} is a simple first order differentiation operator, while on the part in H , it is given by T or T_* , which act essentially like A^* or A with correction or coupling terms from the channels.

6. EXAMPLES

In this section we consider some special cases for which we determine the operator \mathcal{L} and its domain more explicitly.

6.1. The case of bounded imaginary part. We start with a very simple well-known example.

Lemma 6.1. *Assume $A = \Re A + i\Im A$ with $\Re A$ a selfadjoint operator, $\Im A$ a bounded non-negative operator and*

$$D(A) = D(A^*) = D(\Re A). \text{ We can choose } \Gamma = \Gamma_* = \sqrt{2\Im A} \text{ and } E = E_* = \overline{\text{Ran } \Im A} \subseteq H. \text{ Then } U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in$$

$D(\mathcal{L})$ if and only if $u \in D(A)$ and $v_+(0) = v_-(0) + i\Gamma u$.

Proof. It is easy to check (3.1) and (3.2) hold with the given Γ, Γ_*, E and E_* . In particular, Γ and Γ_* are bounded, so for any $v_+(0), v_-(0)$ we have

$$(\Gamma(A + \lambda)^{-1})^* v_+(0) = (A^* + \bar{\lambda})^{-1} \Gamma^* v_+(0) \in D(A^*) \quad \text{and} \quad (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) = (A + \bar{\mu})^{-1} \Gamma_*^* v_-(0) \in D(A).$$

This shows that whenever $U \in D(\mathcal{L})$ we have $u \in D(A)$ and so

$$(6.1) \quad v_+(0) = S^*(-\mu)v_-(0) + i\Gamma(\Gamma_*(A^* + \mu)^{-1})^* v_-(0) + i\Gamma u.$$

Next let $v_-(0)$ in $\text{Ran } \Gamma_*$ be arbitrary, $f \in D(A^*)$ with $v_-(0) = \Gamma_* f$ and $g = (A^* + \mu)f$. Then using the definition of S_* and (3.4), we have

$$\begin{aligned}
(6.2) \quad & S^*(-\mu)v_-(0) + i\Gamma(\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \\
&= [S_*(-\bar{\mu}) + i\Gamma(\Gamma_*(A^* + \mu)^{-1})^*] \Gamma_*(A^* + \mu)^{-1} g \\
&= \Gamma [(A + \bar{\mu})^{-1}(A^* + \bar{\mu})(A^* + \mu)^{-1} - (A + \bar{\mu})^{-1} + (A^* + \mu)^{-1} - (\bar{\mu} - \mu)(A + \bar{\mu})^{-1}] g \\
&= \Gamma f = v_-(0).
\end{aligned}$$

Since this holds for $v_-(0)$ from the dense set $\text{Ran } \Gamma_*$, it holds on the whole space E_* and (6.1) reduces to $v_+(0) = v_-(0) + i\Gamma u$.

On the other hand, if $u \in D(A)$ then clearly the domain inclusions needed in (5.1) are satisfied and (I) follows from $v_+(0) = v_-(0) + i\Gamma u$ by using (6.2). \square

We note that similar considerations work for the case of relatively bounded imaginary part.

6.2. Dissipative Schrödinger operators on the half line. This section considers the combination of dissipative boundary conditions and potentials for Schrödinger operators, providing an example where the imaginary part of the operator is not bounded. It also illustrates the usefulness of being able to consider limits of the parameters λ and μ in the description of the dilation.

We now consider the Schrödinger operator A in $L^2(\mathbb{R}^+)$ as discussed in Example 4.14. We note that the operator A is completely non-selfadjoint provided either $\Im(h) \neq 0$ or $\Im(q)$ is not identically zero. We will prove this in forthcoming work. Our first aim is to determine the asymptotics of the characteristic function given in (4.11). To this end, we begin with an elementary abstract lemma.

Lemma 6.2. *Let $S_n = \begin{pmatrix} \alpha_n & \beta_n \\ \gamma_n & \delta_n \end{pmatrix}$, $n \in \mathbb{N}$, be a sequence of 2×2 block operator matrices in $H_1 \oplus H_2$ consisting of bounded operators and such that S_n is a contraction for all $n \in \mathbb{N}$. If $\alpha_n \xrightarrow{s} I_{H_1}$ and $\delta_n \xrightarrow{s} I_{H_2}$, then $\beta_n, \gamma_n \xrightarrow{s} 0$. Here $A_n \xrightarrow{s} A$ denotes strong convergence.*

Proof. We prove the statement for γ_n . The proof for β_n is similar. For any $x \in H_1$,

$$0 \leq \|\alpha_n^* \alpha_n x - x\|^2 = \|\alpha_n^* \alpha_n x\|^2 + \|x\|^2 - 2 \|\alpha_n x\|^2 \leq 2 \|x\|^2 - 2 \|\alpha_n x\|^2 \rightarrow 0,$$

by assumption on α_n . Thus, $\alpha_n^* \alpha_n \xrightarrow{s} I_{H_1}$. However, $S_n^* S_n$ is also a contraction, so $0 \leq \alpha_n^* \alpha_n + \gamma_n^* \gamma_n \leq I_{H_1}$. This implies $\gamma_n^* \gamma_n \xrightarrow{s} 0$. Then for any $x \in H_1$,

$$\|\gamma_n x\|^2 = \langle \gamma_n^* \gamma_n x, x \rangle \rightarrow 0,$$

proving $\gamma_n \xrightarrow{s} 0$. □

Using that $m_*(z) \sim i\sqrt{z}$ as $z \rightarrow +i\infty$ (see [15]) and using the resolvent estimate for the anti-dissipative operator A^* , we see that the two diagonal terms in the characteristic function in (4.11) converge strongly to the identity. By Lemma 6.2, the two off diagonal terms must converge strongly to 0. Therefore, $S(z) \xrightarrow{s} I$ as $z \rightarrow +i\infty$.

To determine the conditions for lying in $D(\mathcal{L})$ more explicitly, we next determine $(\Gamma_*(A^* - z)^{-1})^*$. Let G_* denote the Green function associated with A^* given by

$$(6.3) \quad G_*(x, y, z) = \frac{1}{\bar{h} - m_*(z)} \begin{cases} (\bar{h}\varphi_*(x, z) + \psi_*(x, z))(m_*(z)\varphi_*(y, z) + \psi_*(y, z)), & x < y, \\ (\bar{h}\varphi_*(y, z) + \psi_*(y, z))(m_*(z)\varphi_*(x, z) + \psi_*(x, z)), & x > y, \end{cases}$$

with φ_* , ψ_* the fundamental solutions from Example 4.14. In particular,

$$(6.4) \quad G_*(0, y, z) = \frac{m_*(z)\varphi_*(y, z) + \psi_*(y, z)}{\bar{h} - m_*(z)}.$$

Let $c \in \mathbb{C}$, $\eta \in L^2(\{\Im q > 0\})$ and $p \in L^2(\mathbb{R}^+)$. Then

$$\begin{aligned} \left\langle (\Gamma_*(A^* - z)^{-1})^* \begin{pmatrix} c \\ \eta(x) \end{pmatrix}, p(x) \right\rangle &= \left\langle \begin{pmatrix} c \\ \eta(x) \end{pmatrix}, \Gamma_*(A^* - z)^{-1} p(x) \right\rangle \\ &= \left\langle \begin{pmatrix} c \\ \eta(x) \end{pmatrix}, \begin{pmatrix} \sqrt{2\Im h} \int_0^\infty G_*(0, y, z) p(y) dy \\ \sqrt{2\Im q(x)} \int_0^\infty G_*(x, y, z) p(y) dy \end{pmatrix} \right\rangle \\ &= \left\langle c\sqrt{2\Im h} \cdot \overline{G_*(0, y, z)} + (A - \bar{z})^{-1} \sqrt{2\Im q(x)} \eta(x), p(y) \right\rangle, \end{aligned}$$

so

$$(\Gamma_*(A^* - z)^{-1})^* \begin{pmatrix} c \\ \eta \end{pmatrix} = c\sqrt{2\Im h} \cdot \overline{G_*(0, \cdot, z)} + (A - \bar{z})^{-1} \sqrt{2\Im q} \eta.$$

The last term clearly lies in $D(A)$, therefore, the condition $u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) \in D(A)$ becomes

$$u + (v_-(0))_1 \sqrt{2\Im h} \cdot \overline{G_*(0, y, z)} = u + (v_-(0))_1 \sqrt{2\Im h} \frac{m_*(z)\varphi_*(y, z) + \psi_*(y, z)}{h - m_*(z)} \in D(A),$$

where $(v_-(0))_1 \in \mathbb{C}$ denotes the first component of $v_-(0)$ and $(v_-(0))_2 \in L^2(\{\Im(q) > 0\})$ denotes the second component of $v_-(0)$. Evaluating the expression and its derivative at 0 gives the condition

$$(6.5) \quad u'(0) - hu(0) = (v_-(0))_1 \sqrt{2\Im h}.$$

Next, we consider the boundary condition $v_+(0) = S^*(-\mu)v_-(0) + i\Gamma(u + (\Gamma_*(A^* + \mu)^{-1})^*v_-(0))$, as $\mu \rightarrow -i\infty$. By $\hat{\Gamma}$ we denote the extension of Γ_* from $D(A^*)$ to $H^1(\mathbb{R}^+)$ by the same formula. From the calculation above we have that

$$(6.6) \quad \hat{\Gamma}(\Gamma_*(A^* + \mu)^{-1})^*v_-(0) = \hat{\Gamma} \left[(v_-(0))_1 \sqrt{2\Im h} \cdot \overline{G_*(0, \cdot, -\mu)} + (A + \bar{\mu})^{-1} \sqrt{2\Im q} (v_-(0))_2 \right].$$

Lemma 6.3. *As $\mu \rightarrow -i\infty$, we have*

$$\hat{\Gamma}(\Gamma_*(A^* + \mu)^{-1})^*v_-(0) \rightarrow 0.$$

Proof. Let $\Im(\mu) < 0$ and recall that $q \in L^\infty(\mathbb{R}^+)$. In a first step we reduce the problem to studying the resolvent of the free operator $(A_0f)(x) = -f''(x)$ with $D(A_0) = D(A)$. Using the Hilbert identity, we have

$$(6.7) \quad \begin{aligned} \hat{\Gamma}(\Gamma_*(A^* + \mu)^{-1})^* &= \hat{\Gamma}(\Gamma_*(A_0^* + \mu)^{-1})^* - \hat{\Gamma}(\Gamma_*(A_0^* + \mu)^{-1}\bar{q}(A^* + \mu)^{-1})^* \\ &= \hat{\Gamma}(\Gamma_*(A_0^* + \mu)^{-1})^* - \hat{\Gamma}[(A + \bar{\mu})^{-1}q(\Gamma_*(A_0^* + \mu)^{-1})^*]. \end{aligned}$$

Now,

$$(6.8) \quad \begin{aligned} \left\| \hat{\Gamma}(A + \bar{\mu})^{-1}q \right\| &\leq \|q\|_\infty \left\| \hat{\Gamma}(A + \bar{\mu})^{-1} \right\| = \|q\|_\infty \left\| \hat{\Gamma}(A_0 + \bar{\mu})^{-1}(A_0 + \bar{\mu})(A + \bar{\mu})^{-1} \right\| \\ &\leq \|q\|_\infty \left\| \hat{\Gamma}(A_0 + \bar{\mu})^{-1} \right\| \|I - q(A + \bar{\mu})^{-1}\| = \|q\|_\infty \left\| \hat{\Gamma}(A_0 + \bar{\mu})^{-1} \right\| \left(1 + \frac{\|q\|_\infty}{|\Im(\mu)|} \right), \end{aligned}$$

using the standard resolvent estimate for the dissipative operator A . Together with (6.7), this means that we are required to show that

$$(a) \left\| \hat{\Gamma}(A_0 + \bar{\mu})^{-1} \right\| \rightarrow 0 \quad \text{and} \quad (b) \left\| \hat{\Gamma}(\Gamma_*(A_0^* + \mu)^{-1})^* \right\| \rightarrow 0, \quad \mu \rightarrow -i\infty.$$

Since, for $u \in H^1(\mathbb{R}^+)$,

$$\hat{\Gamma}u = \begin{pmatrix} \sqrt{2\Im h} u(0) \\ \sqrt{2\Im q} u|_\Omega \end{pmatrix},$$

we have

$$(6.9) \quad \left\| \hat{\Gamma}u \right\| \leq c_h |u(0)| + c_q \|u\|.$$

Letting $P_1u := u(0)$ for $u \in H^1(\mathbb{R}^+)$ and using that taking adjoints preserves the norm, to check (b) above, it is therefore sufficient to show

$$(b)(i) \left\| P_1(\Gamma_*(A_0^* + \mu)^{-1})^* \right\| \rightarrow 0 \quad \text{and} \quad (b)(ii) \left\| \Gamma_*(A_0^* + \mu)^{-1} \right\| \rightarrow 0, \quad \mu \rightarrow -i\infty.$$

From (6.9), since $\Gamma_*v = \hat{\Gamma}v$ for $v \in D(A_0^*)$, we get

$$\begin{aligned} \left\| \Gamma_*(A_0^* + \mu)^{-1}u \right\| &\leq c_h |P_1\Gamma_*(A_0^* + \mu)^{-1}u| + c_q \left\| (A_0^* + \mu)^{-1}u \right\| \\ &\leq c_h |P_1\Gamma_*(A_0^* + \mu)^{-1}u| + \frac{c_q}{|\Im(\mu)|} \|u\|, \end{aligned}$$

using the standard resolvent estimate for the anti-dissipative operator A_0^* . Therefore, (b)(ii) will follow from (b)(i); moreover, (a) will follow from a similar argument. Thus, it remains to prove (b)(i). Replacing A^* by A_0^* and denoting the corresponding Green function by G_0 , from (6.6), we have

$$\left\| P_1(\Gamma_*(A_0^* + \mu)^{-1})^*v_-(0) \right\| \leq c(h, q) \left(|(v_-(0))_1| \cdot |G_0(0, 0, -\mu)| + \|G_0(0, \cdot, -\mu)\| \|(v_-(0))_2\| \right)$$

Now,

$$G_0(0, y, -\mu) = \frac{m_0(-\mu)\varphi_0(y, -\mu) + \psi_0(y, -\mu)}{\bar{h} - m_0(-\mu)},$$

where

$$\varphi_0(y, z) = \frac{e^{i\sqrt{z}y} - e^{-i\sqrt{z}y}}{2i\sqrt{z}}, \quad \psi_0(y, z) = \frac{e^{i\sqrt{z}y} + e^{-i\sqrt{z}y}}{2}$$

and $m_0(z) = i\sqrt{z}$. Therefore, $G_0(0, 0, -\mu) = (\bar{h} - m_0(-\mu))^{-1} \rightarrow 0$ as $\mu \rightarrow -i\infty$ and $|G_0(0, y, -\mu)| \leq (|\bar{h} - m_0(-\mu)|)^{-1} |e^{-\Im(\sqrt{-\mu})y}|$. Hence, $\|G_0(0, y, -\mu)\|_2 \leq (\sqrt{2} |\bar{h} - m_0(-\mu)|)^{-1} \Im(\sqrt{-\mu})^{-1/2} \rightarrow 0$ as $\mu \rightarrow -i\infty$. \square

Combining the lemma with (6.5) we obtain

$$D(\mathcal{L}) = \left\{ U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} : u \in H^2(\mathbb{R}^+), v_{\pm} \in H^1(\mathbb{R}^{\pm}), \right. \\ \left. u'(0) - hu(0) = \sqrt{2\Im h} (v_-(0))_1, v_+(0) = v_-(0) + i\hat{\Gamma}u \right\}.$$

Finally, we consider the action of \mathcal{L} . We have

$$\begin{aligned} T_* \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} &= (A + \bar{\mu})(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) - \bar{\mu}u \\ &= (A + \bar{\mu}) \left(u + \frac{\sqrt{2\Im h}(v_-(0))_1}{h - \overline{m_*(-\mu)}} \overline{(m_*(-\mu)\varphi_*(\cdot, -\mu) + \psi_*(\cdot, -\mu))} + (A + \bar{\mu})^{-1} \sqrt{2\Im q} (v_-(0))_2 \right) - \bar{\mu}u \\ &= -u'' + qu + \sqrt{2\Im q} (v_-(0))_2, \end{aligned}$$

which shows in particular by explicit calculation that T_* here is independent of μ . We now have

$$\mathcal{L}U = \mathcal{L} \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = \begin{pmatrix} v_- \\ -u'' + qu + \sqrt{2\Im q} (v_-(0))_2 \\ iv'_+ \end{pmatrix}.$$

Remark 6.4. *The dilation property of \mathcal{L} can easily be checked: Let $\lambda \in \mathbb{C}^+$ and*

$$\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ w \\ 0 \end{pmatrix}.$$

Then $v_+(x) = v_+(0)e^{-i\lambda x}$, so $v_+(0)$ must vanish. Thus $v_-(0) = -i\hat{\Gamma}u$ and so

$$T_* \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = -u'' + qu + \sqrt{2\Im q} (v_-(0))_2 = -u'' + qu + \sqrt{2\Im q}(-i\sqrt{2\Im q}qu) = -u'' + \bar{q}u,$$

giving $w = -u'' + \bar{q}u$. Therefore, $P_H(\mathcal{L} - \lambda)^{-1}|_H = (A^* - \lambda)^{-1}$. Similarly, one can see that for λ in the lower half plane one has $P_H(\mathcal{L} - \lambda)^{-1}|_H = (A - \lambda)^{-1}$. We will see later in Theorem 7.4 that this is a general property of the operator \mathcal{L} we have constructed.

7. PROPERTIES OF \mathcal{L}

We first calculate the resolvent of the operator \mathcal{L} .

Lemma 7.1. *For $\begin{pmatrix} f \\ w \\ g \end{pmatrix} \in \mathcal{H}$ and $\lambda_0 \in \mathbb{C}^+$, we have*

$$(7.1) \quad (\mathcal{L} - \lambda_0)^{-1} \begin{pmatrix} f \\ w \\ g \end{pmatrix} = \begin{pmatrix} -i\Gamma_*(A^* - \lambda_0)^{-1}we^{-i\lambda_0 x} + iS(\lambda_0) \int_0^\infty e^{i\lambda_0(t-x)}g(t) dt - i \int_0^x e^{i\lambda_0(t-x)}f(t) dt \\ (A^* - \lambda_0)^{-1}w - i(\Gamma(A - \bar{\lambda}_0)^{-1})^* \int_0^\infty e^{i\lambda_0 t}g(t) dt \\ i \int_x^\infty e^{i\lambda_0(t-x)}g(t) dt \end{pmatrix}.$$

Similarly, for $\lambda_0 \in \mathbb{C}^-$, we have

$$(7.2) \quad (\mathcal{L} - \lambda_0)^{-1} \begin{pmatrix} f \\ w \\ g \end{pmatrix} = \begin{pmatrix} -i \int_{-\infty}^x e^{i\lambda_0(t-x)}f(t) dt \\ (A - \lambda_0)^{-1}w + i(\Gamma_*(A^* - \bar{\lambda}_0)^{-1})^* \int_{-\infty}^0 e^{i\lambda_0 t}f(t) dt \\ ie^{i\lambda_0 x}\Gamma(A - \lambda_0)^{-1}w - iS^*(\bar{\lambda}_0) \int_{-\infty}^0 e^{i\lambda_0(t-x)}f(t) dt - i \int_0^x e^{i\lambda_0(t-x)}g(t) dt \end{pmatrix}.$$

Proof. We prove (7.2), the proof of (7.1) is similar. Let $\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in D(\mathcal{L})$, $\lambda_0 \in \mathbb{C}^-$ and

$$(\mathcal{L} - \lambda_0) \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = \begin{pmatrix} iv'_- - \lambda_0 v_- \\ (A + \bar{\mu})(u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0)) - (\bar{\mu} + \lambda_0)u \\ iv'_+ - \lambda_0 v_+ \end{pmatrix} = \begin{pmatrix} f \\ w \\ g \end{pmatrix},$$

where $\mu \in \mathbb{C}^-$ is arbitrary. Choosing $\mu = -\bar{\lambda}_0$, this simplifies to

$$\begin{pmatrix} iv'_- - \lambda_0 v_- \\ (A - \lambda_0)(u + (\Gamma_*(A^* - \bar{\lambda}_0)^{-1})^* v_-(0)) \\ iv'_+ - \lambda_0 v_+ \end{pmatrix} = \begin{pmatrix} f \\ w \\ g \end{pmatrix}.$$

We can easily solve the first and last equation, taking into account that $\Im \lambda_0 < 0$, we get that

$$(7.3) \quad v_-(x) = -i \int_{-\infty}^x e^{i\lambda_0(t-x)} f(t) dt \quad \text{and} \quad v_+(x) = v_+(0) e^{i\lambda_0 x} - i \int_0^x e^{i\lambda_0(t-x)} g(t) dt.$$

Solving the second equation for u gives that

$$u = (A - \lambda_0)^{-1} w - (\Gamma_*(A^* - \bar{\lambda}_0)^{-1})^* v_-(0) = (A - \lambda_0)^{-1} w + i(\Gamma_*(A^* - \bar{\lambda}_0)^{-1})^* \int_{-\infty}^0 e^{i\lambda_0 t} f(t) dt.$$

It remains to determine $v_+(0)$ from

$$\begin{aligned} v_+(0) &= S^*(\bar{\lambda}_0) v_-(0) + i\Gamma(u + (\Gamma_*(A^* - \bar{\lambda}_0)^{-1})^* v_-(0)) \\ &= i\Gamma(A - \lambda_0)^{-1} w - iS^*(\bar{\lambda}_0) \int_{-\infty}^0 e^{i\lambda_0 t} f(t) dt. \end{aligned}$$

Inserting this in (7.3) proves the result. \square

Lemma 7.2. \mathcal{L} is symmetric.

Proof. Let $U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix}, \tilde{U} = \begin{pmatrix} \tilde{v}_- \\ \tilde{u} \\ \tilde{v}_+ \end{pmatrix} \in D(\mathcal{L})$. Then

$$(7.4) \quad \langle \mathcal{L}U, \tilde{U} \rangle - \langle U, \mathcal{L}\tilde{U} \rangle = i(\langle v_-(0), \tilde{v}_-(0) \rangle - \langle v_+(0), \tilde{v}_+(0) \rangle) + \langle T_*U, \tilde{u} \rangle - \langle u, T\tilde{U} \rangle.$$

Then, for $\lambda \in \mathbb{C}^+$ we have

$$\begin{aligned} \langle T_*U, \tilde{u} \rangle - \langle u, T\tilde{U} \rangle &= \langle (A + \lambda)(u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0)) - \lambda u, \tilde{u} \rangle - \langle u, (A^* + \bar{\lambda})(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) - \bar{\lambda} \tilde{u} \rangle \\ &= \langle (A + \lambda)(u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0)), \tilde{u} \rangle - \langle u, (A^* + \bar{\lambda})(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) \rangle \\ &= \langle (A + \lambda)(u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0)), \tilde{u} \rangle \\ &\quad + \langle (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0), (A^* + \bar{\lambda})(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) \rangle \\ &\quad - \langle u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0), (A^* + \bar{\lambda})(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) \rangle \\ &= \langle (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0), (A^* + \bar{\lambda})(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) \rangle \\ &\quad - \langle (A + \lambda)(u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0)), (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0) \rangle \\ &= \langle v_-(0), \Gamma_*(\tilde{u} + (\Gamma(A + \lambda)^{-1})^* \tilde{v}_+(0)) \rangle - \langle \Gamma(u + (\Gamma_*(A^* + \bar{\lambda})^{-1})^* v_-(0)), \tilde{v}_+(0) \rangle. \end{aligned}$$

Using the conditions (I) and (II) in (5.1), we get

$$\begin{aligned} \langle T_*U, \tilde{u} \rangle - \langle u, T\tilde{U} \rangle &= \langle v_-(0), i\tilde{v}_-(0) - iS(-\bar{\lambda})\tilde{v}_+(0) \rangle - \langle -iv_+(0) + iS^*(-\bar{\lambda})v_-(0), \tilde{v}_+(0) \rangle \\ &= \langle v_-(0), i\tilde{v}_-(0) \rangle - \langle -iv_+(0), \tilde{v}_+(0) \rangle. \end{aligned}$$

Together with (7.4), this proves symmetry of \mathcal{L} . \square

Combining the two previous results immediately gives:

Corollary 7.3. \mathcal{L} is selfadjoint.

Proof. By Lemma 7.2, \mathcal{L} is a symmetric operator, while by Lemma 7.1 we have that $\text{Ran}(\mathcal{L} - \lambda) = \mathcal{H}$ for all non-real λ , i.e. \mathcal{L} is a symmetric operator with deficiency indices $(0, 0)$. Hence it is selfadjoint. \square

Theorem 7.4. \mathcal{L} is a minimal selfadjoint dilation of A . In particular, letting $P_H : \mathcal{H} \rightarrow H$ be the projection onto the second component, we have

$$(7.5) \quad P_H(\mathcal{L} - \lambda)^{-1}|_H = \begin{cases} (A - \lambda)^{-1} & \lambda \in \mathbb{C}^-, \\ (A^* - \lambda)^{-1} & \lambda \in \mathbb{C}^+. \end{cases}$$

Proof. The formula (7.5) follows from (7.1) and (7.2) by setting $f = g = 0$ and considering the second component.

It remains to show minimality of the dilation. We need to show that

$$\mathcal{H} = \text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H \\ 0 \end{pmatrix} \right).$$

Let

$$\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \perp \text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H \\ 0 \end{pmatrix} \right).$$

As \mathcal{L} is selfadjoint, $-\lambda(\mathcal{L} - \lambda)^{-1}$ converges strongly to the identity and so $u = 0$. Thus we get that

$$\begin{pmatrix} h_- \\ 0 \\ h_+ \end{pmatrix} := (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} v_- \\ 0 \\ v_+ \end{pmatrix} \perp \begin{pmatrix} 0 \\ H \\ 0 \end{pmatrix}.$$

Then

$$\begin{pmatrix} v_- \\ 0 \\ v_+ \end{pmatrix} = (\mathcal{L} - \lambda) \begin{pmatrix} h_- \\ 0 \\ h_+ \end{pmatrix}, \quad \text{in particular} \quad T \begin{pmatrix} h_- \\ 0 \\ h_+ \end{pmatrix} = 0 = T_* \begin{pmatrix} h_- \\ 0 \\ h_+ \end{pmatrix},$$

which implies from the definition of T and T_* that

$$(7.6) \quad 0 = (A^* + \bar{\lambda})(\Gamma(A + \lambda)^{-1})^* h_+(0), \quad \text{so} \quad (\Gamma(A + \lambda)^{-1})^* h_+(0) = 0 \quad \text{for any} \quad \lambda \in \mathbb{C}^+$$

and

$$(7.7) \quad 0 = (A + \bar{\mu})(\Gamma_*(A^* + \mu)^{-1})^* h_-(0), \quad \text{so} \quad (\Gamma_*(A^* + \mu)^{-1})^* h_-(0) = 0 \quad \text{for any} \quad \mu \in \mathbb{C}^-.$$

From the density of the ranges of $\Gamma : D(A) \rightarrow E$ and $\Gamma_* : D(A^*) \rightarrow E_*$, this implies $h_+(0) = h_-(0) = 0$.

On the other hand, $ih'_+ - \lambda h_+ = v_+$ implies that

$$h_+(t) = e^{-i\lambda t} h_+(0) - ie^{-i\lambda t} \int_0^t e^{i\lambda x} v_+(x) dx.$$

Since for $\lambda \in \mathbb{C}^+$, the function $e^{-i\lambda t}$ is growing, we get for these λ that $h_+(0) = i \int_0^\infty e^{i\lambda x} v_+(x) dx = i\hat{v}_+(\lambda)$. Similarly for $\lambda \in \mathbb{C}^-$, we get that $h_-(0) = -i\hat{v}_-(\lambda)$.

Thus the Fourier transforms of v_+ and v_- vanish in \mathbb{C}^+ and \mathbb{C}^- , respectively, showing that $v_+ = 0 = v_-$, as desired. \square

We complete this section with a discussion of complete non-selfadjointness. We start with a lemma showing independence from parameters of developing certain ranges by the resolvent.

Lemma 7.5. *For any $\lambda', \lambda'' \in \mathbb{C}^+$ we have that*

$$(7.8) \quad \begin{aligned} & \text{clos} \left[\text{Span}_{\mu \in \mathbb{C}^-} (A - \mu)^{-1} \left(\text{Ran} \left((\Gamma(A + i)^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right) \right) \right. \\ & \quad \left. + \text{Span}_{\lambda \in \mathbb{C}^+} (A^* - \lambda)^{-1} \left(\text{Ran} \left((\Gamma(A + i)^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right) \right) \right] \\ & = \text{clos} \left[\text{Span}_{\mu \in \mathbb{C}^-} (A - \mu)^{-1} \left(\text{Ran} \left((\Gamma(A + \lambda')^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - \lambda'')^{-1})^* \right) \right) \right. \\ & \quad \left. + \text{Span}_{\lambda \in \mathbb{C}^+} (A^* - \lambda)^{-1} \left(\text{Ran} \left((\Gamma(A + \lambda')^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - \lambda'')^{-1})^* \right) \right) \right]. \end{aligned}$$

Proof. We will show that the set on the right hand side of (7.8) is contained in the set on the left hand side by considering the difference between two typical terms. Consider

$$\begin{aligned} (A - \mu)^{-1} \left[(\Gamma(A + i)^{-1})^* - (\Gamma(A + \lambda')^{-1})^* \right] &= (A - \mu)^{-1} \left[\Gamma((A + i)^{-1} - (A + \lambda')^{-1}) \right]^* \\ &= (\bar{\lambda}' + i)(A - \mu)^{-1} \left[\Gamma((A + i)^{-1}(A + \lambda')^{-1}) \right]^* \\ &= (\bar{\lambda}' + i)(A - \mu)^{-1}(A^* + \bar{\lambda}')^{-1} \left[\Gamma(A + i)^{-1} \right]^*, \end{aligned}$$

where we have used the Hilbert identity. Next, we use (3.4) to obtain

$$(7.9) \quad \begin{aligned} & (A - \mu)^{-1} \left[(\Gamma(A + i)^{-1})^* - (\Gamma(A + \lambda')^{-1})^* \right] \\ &= \frac{\bar{\lambda}' + i}{\bar{\lambda}' + \mu} \left[(A - \mu)^{-1} - (A^* + \bar{\lambda}')^{-1} + i(\Gamma_*(A^* - \bar{\mu})^{-1})^* (\Gamma_*(A^* + \bar{\lambda}')^{-1}) \right] \left[\Gamma(A + i)^{-1} \right]^*. \end{aligned}$$

Clearly, the first two terms on the right lie in the desired set. For the last term, we note the following two facts: Since $-\mu(A - \mu)^{-1} \rightarrow I$ as $\mu \rightarrow -i\infty$, the set on the left hand side of (7.8) contains $\text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right)$ and by the Hilbert identity we have

$$(A - \mu)^{-1} (\Gamma_*(A^* - i)^{-1})^* = (\Gamma_*(A^* - i)^{-1}(A^* - \bar{\mu})^{-1})^* = \frac{1}{i - \bar{\mu}} \left[(\Gamma_*(A^* - i)^{-1})^* - (\Gamma_*(A^* - \bar{\mu})^{-1})^* \right].$$

Thus $\text{Ran} \left((\Gamma_*(A^* - \bar{\mu})^{-1})^* \right) \subseteq \text{Span} \left[\left(\text{Ran} \left((A - \mu)^{-1} (\Gamma_*(A^* - i)^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right) \right) \right]$, showing that all terms on the right hand side of (7.9) lie in the set on the left hand side of (7.8).

All other inclusions of terms on the right of (7.8) in the set on the left can be checked similarly. The reverse inclusion follows in a similar manner. \square

In the following we present a construction of the Langer decomposition from Proposition 2.10 and show its relation to the dilation.

Theorem 7.6. *Let A be a maximal dissipative operator and denote*

$$(7.10) \quad \begin{aligned} H_{cns} = & \text{clos} \left[\text{Span}_{\mu \in \mathbb{C}^-} \left\{ (A - \mu)^{-1} \left(\text{Ran} \left((\Gamma(A + i)^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right) \right) \right\} \right. \\ & \left. + \text{Span}_{\lambda \in \mathbb{C}^+} \left\{ (A^* - \lambda)^{-1} \left(\text{Ran} \left((\Gamma(A + i)^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - i)^{-1})^* \right) \right) \right\} \right]. \end{aligned}$$

Let $H_{sa} = H \ominus H_{cns}$. Then

- (1) H_{cns} is a reducing subspace for A .
- (2) $A_{sa} := A|_{H_{sa}}$ is selfadjoint.
- (3) The operator $A_{cns} := A|_{H_{cns}}$ is completely non-selfadjoint.
- (4) The subspace $\begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix}$ is reducing for the dilation \mathcal{L} with \mathcal{L} restricted to $\begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix}$ given by A_{sa} .
- (5) We have

$$\text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \right) = \mathcal{H} \ominus \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} = \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ H_{cns} \\ L^2(\mathbb{R}^+, E) \end{pmatrix}.$$

- (6) \mathcal{L} restricted to $\begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ H_{cns} \\ L^2(\mathbb{R}^+, E) \end{pmatrix}$ is a minimal selfadjoint dilation of A_{cns} .

Remark 7.7. *For the case of bounded imaginary part, this result is known and can be found in [30].*

Proof. (1) We show that $(A - \mu_0)^{-1} H_{cns} \subseteq H_{cns}$ for all $\mu_0 \in \mathbb{C}^-$. Similarly, one can show that $(A^* - \lambda_0)^{-1} H_{cns} \subseteq H_{cns}$ for all $\lambda_0 \in \mathbb{C}^+$. Together, this shows that H_{cns} is reducing for the resolvent of A , which implies it is reducing for A .

We consider $w \in H_{cns}$ of the form

$$\begin{aligned} w = & (A - \mu_1)^{-1} (\Gamma(A + i)^{-1})^* h_1 + (A - \mu_2)^{-1} (\Gamma_*(A^* - i)^{-1})^* h_2^* \\ & + (A^* - \lambda_2)^{-1} (\Gamma(A + i)^{-1})^* h_2 + (A^* - \lambda_1)^{-1} (\Gamma_*(A^* - i)^{-1})^* h_1^* =: w_1 + w_2 + w_3 + w_4, \end{aligned}$$

where $\mu_1, \mu_2 \in \mathbb{C}^-$, $\lambda_1, \lambda_2 \in \mathbb{C}^+$, $h_1, h_2 \in E$ and $h_1^*, h_2^* \in E_*$. It is sufficient to show $(A - \mu_0)^{-1} w_i \in H_{cns}$ for $i = 1, \dots, 4$, since linear combinations of vectors of this form are dense and the resolvent bounded.

It is immediately clear from the Hilbert identity that $(A - \mu_0)^{-1} w_1, (A - \mu_0)^{-1} w_2 \in H_{cns}$. For $(A - \mu_0)^{-1} w_3, (A - \mu_0)^{-1} w_4$, we use (3.4) and the fact that $-\mu(A - \mu)^{-1} \rightarrow I$ as $\mu \rightarrow -i\infty$ to prove the inclusion.

- (2) Let $u \in H_{sa}$. By Lemma 7.5 and again using that $-\mu(A - \mu)^{-1} \rightarrow I$ as $\mu \rightarrow -i\infty$, this implies that for any $\mu \in \mathbb{C}^-$ we have

$$u \perp \text{Ran} \left((\Gamma_*(A^* + \mu)^{-1})^* \right), \text{ or equivalently, } \Gamma_*(A^* + \mu)^{-1} u = 0.$$

By a similar argument, $\Gamma(A + \lambda)^{-1} u = 0$ for any $\lambda \in \mathbb{C}^+$. Using (3.3) and (3.4), respectively, we get that

$$\left[(A + \lambda)^{-1} - (A^* + \mu)^{-1} + (\lambda - \mu)(A^* + \mu)^{-1}(A + \lambda)^{-1} \right] u = 0$$

and

$$\left[(A + \lambda)^{-1} - (A^* + \mu)^{-1} + (\lambda - \mu)(A + \lambda)^{-1}(A^* + \mu)^{-1} \right] u = 0.$$

Choosing $\lambda = \bar{\mu}$, we get

$$(7.11) \quad [(A + \bar{\mu})^{-1} - (A^* + \mu)^{-1} + (\bar{\mu} - \mu)(A^* + \mu)^{-1}(A + \bar{\mu})^{-1}]|_{H_{sa}} = 0$$

and

$$(7.12) \quad [(A + \bar{\mu})^{-1} - (A^* + \mu)^{-1} + (\bar{\mu} - \mu)(A + \bar{\mu})^{-1}(A^* + \mu)^{-1}]|_{H_{sa}} = 0.$$

Next, let T be the Cayley transform of A_{sa} at μ , i.e. $T = (A_{sa} + \mu)(A_{sa} + \bar{\mu})^{-1} = I - (\bar{\mu} - \mu)(A_{sa} + \bar{\mu})^{-1}$. Then $(A_{sa} + \bar{\mu})^{-1} = (\mu - \bar{\mu})^{-1}(T - I)$, $(A_{sa}^* + \mu)^{-1} = (\bar{\mu} - \mu)^{-1}(T^* - I)$ and (7.11) is equivalent to

$$0 = T - I + T^* - I + (T^* - I)(T - I) = T + T^* - 2I + T^*T - T^* - T + I = T^*T - I,$$

so $T^*T = I$. Similarly, (7.12) is equivalent to $TT^* = I$, so T is unitary and its inverse Cayley transform A_{sa} is selfadjoint.

(3) Assume $W \subseteq H_{cns}$ is a reducing subspace such that $A|_W$ is selfadjoint. From (3.3) and (3.4), we get that

$$(7.13) \quad [(A + i)^{-1} - (A^* - i)^{-1} + 2i(A^* - i)^{-1}(A + i)^{-1}]|_W = -i(\Gamma(A + i)^{-1})^*(\Gamma(A + i)^{-1})|_W$$

and

$$(7.14) \quad [(A + i)^{-1} - (A^* - i)^{-1} + 2i(A + i)^{-1}(A^* - i)^{-1}]|_W = -i(\Gamma_*(A^* - i)^{-1})^*(\Gamma_*(A^* - i)^{-1})|_W.$$

As we are assuming that $A|_W$ is selfadjoint, the left hand sides of (7.13) and (7.14) vanish. Due to the density of the ranges of Γ and Γ_* , this implies that $\Gamma(A + i)^{-1}|_W = \Gamma_*(A^* - i)^{-1}|_W = 0$. By the same reasoning, we get that $\Gamma(A + \lambda)^{-1}|_W = 0$ for any $\lambda \in \mathbb{C}^+$ and $\Gamma_*(A^* + \mu)^{-1}|_W = 0$ for any $\mu \in \mathbb{C}^-$. Thus $u \in W$ implies that

$$u \perp \text{Ran} (\Gamma(A + \lambda)^{-1})^* \quad \text{and} \quad u \perp \text{Ran} (\Gamma_*(A^* + \mu)^{-1})^*.$$

As W is reducing, this also implies that for any $\lambda' \in \mathbb{C}^+$

$$(A + \lambda')^{-1}u \perp \text{Ran} (\Gamma(A + \lambda)^{-1})^* \quad \text{and} \quad (A + \lambda')^{-1}u \perp \text{Ran} (\Gamma_*(A^* + \mu)^{-1})^*,$$

which shows that

$$u \perp (A^* + \bar{\lambda}')^{-1} \left(\text{Ran} (\Gamma(A + \lambda)^{-1})^* \bigcup \text{Ran} (\Gamma_*(A^* + \mu)^{-1})^* \right).$$

By the same reasoning for any $\mu' \in \mathbb{C}^-$ we have

$$u \perp (A + \bar{\mu}')^{-1} \left(\text{Ran} (\Gamma(A + \lambda)^{-1})^* \bigcup \text{Ran} (\Gamma_*(A^* + \mu)^{-1})^* \right),$$

so $u \perp H_{cns}$ and $u = 0$.

(4) We show that the subspace $\begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix}$ is reducing for the resolvent of \mathcal{L} . Let $w \in H_{sa}$. Since $w \perp H_{cns}$,

we know from the Hilbert identity that for any $\lambda \in \mathbb{C}^-$ we have $\Gamma(A - \lambda)^{-1}w = 0 = \Gamma_*(A^* - \bar{\lambda})^{-1}w$. Thus from (7.1) and (7.2), we get

$$(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ w \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ (A^* - \lambda)^{-1}w \\ 0 \end{pmatrix} \quad \text{and} \quad (\mathcal{L} - \bar{\lambda})^{-1} \begin{pmatrix} 0 \\ w \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ (A - \bar{\lambda})^{-1}w \\ 0 \end{pmatrix}.$$

The claim now follows immediately from part (1).

(5) Suppose

$$\begin{pmatrix} f \\ w \\ g \end{pmatrix} \perp (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \quad \text{for all } \lambda \notin \mathbb{R}.$$

Since \mathcal{L} is symmetric, this means that

$$(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} f \\ w \\ g \end{pmatrix} \perp \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \quad \text{for all } \lambda \notin \mathbb{R}.$$

First, let $\lambda \in \mathbb{C}^+$. Then by (7.1), we see that

$$0 = \int_x^\infty e^{i\lambda(t-x)}g(t) dt = \int_0^\infty e^{i\lambda t}g(t) dt,$$

which implies $g \equiv 0$. Similarly, choosing $\lambda \in \mathbb{C}^-$, by (7.2), we see that $f \equiv 0$.

From the first component in (7.1), we now see that $\Gamma_*(A^* - \lambda)^{-1}w = 0$ and from the third component of (7.2), we have $\Gamma(A - \bar{\lambda})^{-1}w = 0$ for all $\lambda \in \mathbb{C}^+$. Thus

$$w \perp \text{Ran} \left((\Gamma(A - \bar{\lambda})^{-1})^* \right) \cup \text{Ran} \left((\Gamma_*(A^* - \lambda)^{-1})^* \right).$$

Since \mathcal{L} is selfadjoint, it follows immediately from the Hilbert identity that

$$\text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \right)$$

is a reducing subspace and therefore, for any $\mu, \lambda \notin \mathbb{R}$ also

$$(\mathcal{L} - \mu)^{-1}(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ w \\ 0 \end{pmatrix} \perp \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix}.$$

Now, choosing $\mu, \lambda \in \mathbb{C}^+$, from (7.2), there exists \tilde{g} such that

$$(\mathcal{L} - \mu)^{-1}(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ w \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ (A - \bar{\lambda})^{-1}w \\ \tilde{g} \end{pmatrix}$$

and by repeating the arguments above, we see that $\Gamma_*(A^* - \mu)^{-1}(A - \bar{\lambda})^{-1}w = 0$, giving that

$$w \perp (A^* - \lambda)^{-1} \text{Ran} \left((\Gamma_*(A^* - \mu)^{-1})^* \right).$$

Similarly, choosing μ and λ from appropriate half-planes we see that $w \perp H_{cns}$. Therefore, we have shown that

$$(7.15) \quad \left(\text{clos} \left[\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \right] \right)^\perp \subseteq \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix}.$$

On the other hand, using part (4), we know that

$$(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} \subseteq \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix}.$$

Taking orthogonal complements, this gives that

$$\begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \subseteq \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ H_{cns} \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \subseteq \left[(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} \right]^\perp$$

and so

$$(\mathcal{L} - \lambda)^{-1} \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ 0 \\ L^2(\mathbb{R}^+, E) \end{pmatrix} \subseteq \left[\begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} \right]^\perp.$$

Taking the linear span, this together with (7.15) gives (5).

- (6) Since we have shown (4) and \mathcal{L} is selfadjoint, it is clear that \mathcal{L} restricted to $\begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ H_{cns} \\ L^2(\mathbb{R}^+, E) \end{pmatrix}$ is a selfadjoint dilation of A_{cns} . Since \mathcal{L} is a minimal dilation of A by Theorem 7.4, we get

$$(7.16) \quad \mathcal{H} = \text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \left[\begin{pmatrix} 0 \\ H_{cns} \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} \right] \right)$$

$$(7.17) \quad \subseteq \text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H_{cns} \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ H_{sa} \\ 0 \end{pmatrix} \right),$$

using part (4). This proves that

$$\text{clos} \left(\text{Span}_{\lambda \notin \mathbb{R}} (\mathcal{L} - \lambda)^{-1} \begin{pmatrix} 0 \\ H_{cns} \\ 0 \end{pmatrix} \right) = \begin{pmatrix} L^2(\mathbb{R}^-, E_*) \\ H_{cns} \\ L^2(\mathbb{R}^+, E) \end{pmatrix}$$

and hence the required minimality. □

The next result gives several descriptions of a core for \mathcal{L} .

Theorem 7.8. (1) *The set*

$$\mathcal{C} := \left\{ \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in D(\mathcal{L}) : v_+(0) \in \text{Ran } \Gamma \right\}$$

is a core for \mathcal{L} , i.e. it is dense in $D(\mathcal{L})$ in the graph norm.

(2) *We have the following equivalent descriptions of \mathcal{C} :*

$$(7.18) \quad \mathcal{C} = \left\{ \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in D(\mathcal{L}) : v_-(0) \in \text{Ran } \Gamma_* \right\}$$

$$(7.19) \quad = \left\{ \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in \mathcal{H} : v_{\pm} \in H^1, \exists h \in D(A). h - iu \in D(A^*), v_+(0) = \Gamma h, v_-(0) = \Gamma_*(h - iu) \right\}$$

$$(7.20) \quad = \left\{ \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in \mathcal{H} : v_{\pm} \in H^1, \exists h_* \in D(A^*). h_* + iu \in D(A), v_-(0) = \Gamma_* h_*, v_+(0) = \Gamma(h_* + iu) \right\}$$

(3) *For $\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in \mathcal{C}$ we have*

$$(7.21) \quad \mathcal{L} \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} = \begin{pmatrix} i v'_- \\ i [A^*(h - iu) - Ah] \\ i v'_+ \end{pmatrix},$$

where $h \in D(A)$ is as in (7.19).

Proof. (1) Let \perp_G denote orthogonality in the graph norm and $W = \begin{pmatrix} f \\ w \\ g \end{pmatrix} \in \mathcal{C}^{\perp_G} \subseteq D(\mathcal{L})$. Choosing $U =$

$\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in D(\mathcal{L})$ with $v_- = 0$, $u = 0$ and $v_+ \in H_0^1(\mathbb{R}^+, E)$ arbitrary, the orthogonality condition gives

$$0 = \langle g, v_+ \rangle + \langle g', v'_+ \rangle.$$

This implies that $g' \in H^1(\mathbb{R}^+, E)$ and that $g - g'' = 0$ in the sense of distributions, so $g(x) = g(0)e^{-x}$. Similarly, choosing $v_+ = 0$, $u = 0$ and $v_- \in H_0^1(\mathbb{R}^-, E_*)$ arbitrary shows that $f(x) = f(0)e^x$.

Next, choose $v_- = 0$. Then by (I) in (5.1) we have $v_+(0) = i\Gamma u$ and the orthogonality condition yields

$$\begin{aligned} 0 &= \langle w, u \rangle + \langle TW, T_*U \rangle + \langle g, v_+ \rangle + \langle g', v'_+ \rangle \\ &= \langle w, u \rangle + \langle (A^* + \bar{\lambda})(w + (\Gamma(A + \lambda)^{-1})^*g(0)) - \bar{\lambda}w, Au \rangle + \langle g(0), i\Gamma u \rangle. \end{aligned}$$

Since $u \in D(A)$ we may set $u = (A + i)^{-1}b$ for some $b \in H$. Then

$$0 = \left\langle (A^* - i)^{-1}w + (A(A + i)^{-1})^* [(A^* + \bar{\lambda})(w + (\Gamma(A + \lambda)^{-1})^*g(0)) - \bar{\lambda}w] - i(\Gamma(A + i)^{-1})^*g(0), b \right\rangle.$$

Since $b \in H$ is arbitrary and choosing $\lambda = i$, we get

$$0 = (A^* - i)^{-1}w + (A(A + i)^{-1})^* [(A^* - i)(w + (\Gamma(A + i)^{-1})^*g(0)) + iw] - i(\Gamma(A + i)^{-1})^*g(0).$$

Next use that

$$(A(A + i)^{-1})^* = (I - i(A + i)^{-1})^* = I + i(A^* - i)^{-1},$$

to get

$$\begin{aligned} 0 &= (A^* - i)^{-1}w + (A^* - i)(w + (\Gamma(A + i)^{-1})^*g(0)) + iw \\ &\quad + i(A^* - i)^{-1}[(A^* - i)(w + (\Gamma(A + i)^{-1})^*g(0)) + iw] - i(\Gamma(A + i)^{-1})^*g(0) \\ &= (A^* - i)(w + (\Gamma(A + i)^{-1})^*g(0)) + 2iw. \end{aligned}$$

Therefore, applying $(A^* - i)^{-1}$, we get

$$(7.22) \quad (I + 2i(A^* - i)^{-1})w = -(\Gamma(A + i)^{-1})^*g(0).$$

We now choose U with $v_+ = 0$, $u = (A^* - i)^{-1}c$ and $v_-(0) = -i\Gamma_*u$. Then

$$\begin{aligned} 0 &= \langle w, u \rangle + \langle T_*W, TU \rangle + \langle f, v_- \rangle + \langle f', v'_- \rangle \\ &= \langle w, u \rangle + \langle (A + i)(w + (\Gamma_*(A^* - i)^{-1})^*f(0)) - iw, A^*u \rangle + \langle f(0), -i\Gamma_*u \rangle. \end{aligned}$$

A similar calculation to above shows that

$$(7.23) \quad (I - 2i(A + i)^{-1})w = -(\Gamma_*(A^* - i)^{-1})^*f(0).$$

We now go back to the case when $v_- = 0$ and $u = (A + i)^{-1}b$ and write the orthogonality relation using the expression for T_* rather than T for W , i.e.

$$\begin{aligned} 0 &= \langle w, u \rangle + \langle T_*W, T_*U \rangle + \langle g, v_+ \rangle + \langle g', v'_+ \rangle \\ &= \langle w, u \rangle + \langle (A + i)(w + (\Gamma_*(A^* - i)^{-1})^*f(0)) - iw, Au \rangle + \langle g(0), i\Gamma u \rangle. \end{aligned}$$

Using (7.23) and (7.22), this gives

$$\begin{aligned} 0 &= \langle w, u \rangle + \langle iw, Au \rangle + \langle g(0), i\Gamma u \rangle \\ &= \langle (A^* - i)^{-1}w + i(I + i(A^* - i)^{-1})w + i(I + 2i(A^* - i)^{-1})w, b \rangle \\ &= \langle 2iw - 2(A^* - i)^{-1}w, b \rangle. \end{aligned}$$

Thus $(I + i(A^* - i)^{-1})w = 0$, implying $w \in D(A^*)$. Applying $A^* - i$, we then find $w \in \ker(A^*)$. This implies $w \in H_{sa}$.

On the other hand, from (7.22), we now get

$$-w = (I + 2i(A^* - i)^{-1})w = -(\Gamma(A + i)^{-1})^*g(0) \in H_{cns}.$$

Hence, $w \in H_{sa} \cap H_{cns} = \{0\}$.

Equations (7.22) and (7.23) now give

$$(7.24) \quad (\Gamma_*(A^* - i)^{-1})^*f(0) = 0 = (\Gamma(A + i)^{-1})^*g(0).$$

Then (I) in (5.1) with $\mu = -i$ and using (7.24) simply becomes $g(0) = S^*(i)f(0)$ and for any $h \in D(A)$ we have

$$\begin{aligned} \langle g(0), \Gamma h \rangle &= \langle f(0), S(i)\Gamma h \rangle = \langle f(0), \Gamma_*(A^* - i)^{-1}(A - i)h \rangle \\ &= \langle (\Gamma_*(A^* - i)^{-1})^*f(0), (A - i)h \rangle = 0, \end{aligned}$$

where we have again used (7.24). Hence, $g(0)$ is orthogonal to $\text{Ran } \Gamma$, which is dense in E , so $g(0) = 0$. From (II) in (5.1) we get $f(0) = 0$, which completes the proof of the core property.

(2) If $v_+(0) \in \text{Ran } \Gamma$, then from (II) in (5.1), we see that $v_-(0) \in \text{Ran } \Gamma_*$. The converse follows from condition (I) in $D(\mathcal{L})$. Thus (7.18) holds.

To show (7.19), first let $U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in \mathcal{C}$ and let $w_+ \in H$ be such that $v_+(0) = \Gamma(A + \lambda)^{-1}w_+$ for some

$\lambda \in \mathbb{C}^+$. Using (3.3), we have

$$(7.25) \quad \begin{aligned} \mu + (\Gamma(A + \lambda)^{-1})^*v_+(0) &= u + (\Gamma(A + \lambda)^{-1})^*\Gamma(A + \lambda)^{-1}w_+ \\ &= u + i((A + \lambda)^{-1}w_+ - (A^* + \bar{\lambda})^{-1}w_+ + (\lambda - \bar{\lambda})(A^* + \bar{\lambda})^{-1}(A + \lambda)^{-1}w_+), \end{aligned}$$

Thus, from the domain condition in (5.1), we get that $u + i(A + \lambda)^{-1}w_+ \in D(A^*)$. Set $h = (A + \lambda)^{-1}w_+ \in D(A)$. Clearly, $h - iu \in D(A^*)$ and $v_+(0) = \Gamma h$. It remains to show that $v_-(0) = \Gamma_*(h - iu)$. By the previous calculation, condition (II) in (5.1) now becomes

$$(7.26) \quad \begin{aligned} v_-(0) &= S(-\bar{\lambda})\Gamma h - i\Gamma_*(u + (\Gamma(A + \lambda)^{-1})^*\Gamma(A + \lambda)^{-1}w_+) \\ &= S(-\bar{\lambda})\Gamma h + \Gamma_*[-iu + h - (A^* + \bar{\lambda})^{-1}w_+ + (\lambda - \bar{\lambda})(A^* + \bar{\lambda})^{-1}(A + \lambda)^{-1}w_+] \\ &= \Gamma_*(h - iu), \end{aligned}$$

using the definition of S , (4.3).

On the other hand, if $U = \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix}$ lies in the set on the r.h.s. of (7.19), then clearly $v_-(0) \in \text{Ran } \Gamma_*$ and it remains to check that $u + (\Gamma(A + \lambda)^{-1})^* v_+(0) \in D(A^*)$ and (II) holds. Setting $v_+(0) = \Gamma h = \Gamma(A + \lambda)^{-1} w_+$ this follows by the same calculations as in (7.25) and (7.26). This proves (7.19).

Finally, (7.20) follows by setting $h_* = h - iu$.

(3) Let $\begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} \in \mathcal{C}$. Then $v_+(0) = \Gamma h = \Gamma(A + \lambda)^{-1} v$ with h as in (7.19) and some $v \in H$. Therefore, using (5.9) we have

$$\begin{aligned} T \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} &= (A^* + \mu) (u + (\Gamma(A + \bar{\mu})^{-1})^* v_+(0)) - \mu u \\ &= (A^* + \mu) (u + (\Gamma(A + \bar{\mu})^{-1})^* \Gamma(A + \lambda)^{-1} v) - \mu u. \end{aligned}$$

Using (3.3), this gives

$$\begin{aligned} T \begin{pmatrix} v_- \\ u \\ v_+ \end{pmatrix} &= (A^* + \mu) (u + i((A + \lambda)^{-1} - (A^* + \mu)^{-1} + (\lambda - \mu)(A^* + \mu)^{-1}(A + \lambda)^{-1})v) - \mu u \\ &= (A^* + \mu) (u + ih) - iv + i(\lambda - \mu)(A + \lambda)^{-1}v - \mu u \\ &= (A^* + \mu) (u + ih) - i(A + \lambda)h + i(\lambda - \mu)h - \mu u \\ &= iA^*(h - iu) - iAh. \end{aligned}$$

The statement now follows from the definition on \mathcal{L} . \square

8. DISCUSSION

8.1. Advantages of our Construction. We compare the construction of the operators Γ, Γ_* in our model to having to determine the square root of operators in other models. We consider the case when the imaginary part has finite rank: $A = \Re A + iV$ with V of finite rank. Then we need to determine Γ so that

$$2 \langle Vu, v \rangle = \langle \Gamma u, \Gamma v \rangle.$$

V can be represented by a positive Hermitian matrix. Using the Cholesky decomposition, we can write $2V = \Gamma^* \Gamma$ for an upper triangular matrix Γ with non-negative diagonal entries. Therefore, our method requires calculating the Cholesky decomposition of the matrix rather than its square root.

8.2. Comparison to the Kudryashov/Ryzhov model. Based on the work of Kudryashov, in [40], Ryzhov discusses two selfadjoint dilations (which are then shown to coincide) of a dissipative operator A . These are constructed using the Sz.-Nagy-Foias functional model involving square roots, as discussed at the end of Section 4.

We show that for the special choice of $\lambda = i$, our model can be recovered from the results in [40]. However, the method will not reproduce our explicit formulae, as transformations that use square roots of operators are involved.

Let A be a maximal dissipative operator, T its Cayley transform (2.1) and define D_T and D_{T^*} as in (4.14) and (4.15). Set $\tilde{E} = \overline{\text{Ran } D_T} \subseteq H$, $\tilde{E}_* = \overline{\text{Ran } D_{T^*}} \subseteq H$, $Q = D_T/\sqrt{2}$, $Q_* = D_{T^*}/\sqrt{2}$, let $\tilde{\Gamma} = Q(A + i) : D(A) \rightarrow H$, $\tilde{\Gamma}_* = Q_*(A^* - i) : D(A^*) \rightarrow H$. The selfadjoint dilation \mathcal{A} of A in [40] is then given by

$$(8.1) \quad D(\mathcal{A}) = \left\{ \begin{pmatrix} h_- \\ h_0 \\ h_+ \end{pmatrix} : h_0 \in H, h_+ \in H^1(\mathbb{R}^+, \tilde{E}), h_- \in H^1(\mathbb{R}^-, \tilde{E}_*), \right. \\ \left. h_0 + Q_* h_-(0) \in D(A), h_+(0) = T^* h_-(0) + i\tilde{\Gamma} (h_0 + Q_* h_-(0)) \right\}$$

acting in the space $L^2(\mathbb{R}^-, \tilde{E}_*) \oplus H \oplus L^2(\mathbb{R}^+, \tilde{E})$ with

$$\mathcal{A} \begin{pmatrix} h_- \\ h_0 \\ h_+ \end{pmatrix} = \begin{pmatrix} ih'_- \\ -ih_0 + (A + i)(h_0 + Q_* h_-(0)) \\ ih'_+ \end{pmatrix} \quad \text{for} \quad \begin{pmatrix} h_- \\ h_0 \\ h_+ \end{pmatrix} \in D(\mathcal{A}).$$

By Lemma 3.3, there exist unitary operators $U : \tilde{E} \rightarrow E$ and $U_* : \tilde{E}_* \rightarrow E_*$ such that $\Gamma = U\tilde{\Gamma}$ and $\Gamma_* = U_*\tilde{\Gamma}_*$. Here, E, E_*, Γ and Γ_* are as in Lemma 3.1.

Lemma 8.1. *We have that*

$$\begin{pmatrix} U_* & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & U \end{pmatrix} \mathcal{A} = \mathcal{L} \begin{pmatrix} U_* & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & U \end{pmatrix}.$$

Proof. We first show that $h_0 + Q_* h_-(0) \in D(A)$ if and only if $u + (\Gamma_*(A^* - i)^{-1})^* v_-(0) \in D(A)$, where $u = h_0$ and $v_-(0) = U_* h_-(0)$. Recalling that Q_* is selfadjoint, we get

$$(8.2) \quad \begin{aligned} u + (\Gamma_*(A^* + \mu)^{-1})^* v_-(0) &= h_0 + (U_* \tilde{\Gamma}_*(A^* - i)^{-1})^* U_* h_-(0) \\ &= h_0 + (U_* Q_*)^* U_* h_-(0) = h_0 + Q_* h_-(0). \end{aligned}$$

Next we show that $h_+(0) = T^* h_-(0) + i\tilde{\Gamma}(h_0 + Q_* h_-(0))$ is equivalent to condition (I) in (5.1). By (8.2), condition (I) is equivalent to

$$U h_+(0) = S^*(i) U_* h_-(0) + i U \tilde{\Gamma}(h_0 + Q_* h_-(0)).$$

We therefore need to show that $U^* S^*(i) U_* = T^*$ or, equivalently, $S(i) = U_* T U^*$. By (4.3), on $D(A)$ we have

$$\begin{aligned} S(i) U \tilde{\Gamma} &= S(i) \Gamma = \Gamma_*(A^* - i)^{-1} (A - i) = U_* \tilde{\Gamma}_*(A^* - i)^{-1} (A - i) = U_* Q_*(A - i) \\ &= U_* Q_*(A - i)(A + i)^{-1} (A + i) = U_* Q_* T (A + i) = U_* T Q (A + i) = U_* T \tilde{\Gamma}, \end{aligned}$$

where we have used that $Q_* T = T Q$. Thus on the dense set $\text{Ran } \tilde{\Gamma}$, we have $S(i) U = U_* T$, as required.

It remains to show that the action of the operators coincides. Clearly, the action on the incoming and outgoing channels does. That the action on the middle component coincides, follows from (8.2) together with (5.10), Lemma 5.7 and Corollary 5.8. \square

8.3. Connection of the M -function to the characteristic function in the case of a symmetric minimal operator. In [39], Ryzhov develops a functional model for certain non-selfadjoint extensions of a symmetric operator with equal deficiency indices by using the classical boundary triple framework (see [7, 11, 28, 39] for the definition and details). We now compare the M -function $M(\lambda)$ arising in the boundary triple framework to our characteristic function $S(\lambda)$ in the case of an underlying symmetric operator. Related results and connections to scattering theory can be found in [3, 4], while [2] gives another construction of a minimal selfadjoint dilation in the case on an underlying symmetric operator. We stress that in our construction neither symmetry of the underlying operator nor equal deficiency indices are required.

Let L be a symmetric operator. We construct an associated boundary triple using the von Neumann formula: $D(L^*) = D(L) \dot{+} N_+ \dot{+} N_-$, where $N_{\pm} = \ker(L^* \mp i)$. Let $f = f_0 + f_i + f_{-i}$ and $g = g_0 + g_i + g_{-i}$ lie in $D(L^*)$ and be decomposed according to the von Neumann formula. Then

$$\langle L^* f, g \rangle - \langle f, L^* g \rangle = 2i (\langle f_i, g_i \rangle - \langle f_{-i}, g_{-i} \rangle) = \langle \Gamma_1 f, \tilde{\Gamma}_0 g \rangle - \langle \Gamma_0 f, \tilde{\Gamma}_1 g \rangle,$$

where $\Gamma_1 f = \sqrt{2} f_i$, $\Gamma_0 f = \sqrt{2} f_{-i}$, $\tilde{\Gamma}_0 g = -i\sqrt{2} g_i$ and $\tilde{\Gamma}_1 g = -i\sqrt{2} g_{-i}$. We also have

$$(8.3) \quad \langle L^* f, g \rangle - \langle f, L^* g \rangle = i \langle \Gamma_1 f, \Gamma_1 g \rangle_{N_+} - i \langle \Gamma_0 f, \Gamma_0 g \rangle_{N_-}.$$

Moreover, for $B : N_- \rightarrow N_+$, let $L_B = L^*|_{D(L_B)}$, where $D(L_B) = \ker(\Gamma_1 - B\Gamma_0)$. Then for $f, g \in D(L_B)$,

$$(8.4) \quad \langle L_B f, g \rangle - \langle f, L_B g \rangle = i \langle B\Gamma_0 f, B\Gamma_0 g \rangle_{N_+} - i \langle \Gamma_0 f, \Gamma_0 g \rangle_{N_-},$$

so L_B is dissipative if and only if $B^* B \geq I_{N_-}$. For $f, g \in D(L_B^*) = \ker(\tilde{\Gamma}_1 - B^* \tilde{\Gamma}_0)$, we have

$$(8.5) \quad \langle L_B^* f, g \rangle - \langle f, L_B^* g \rangle = i \langle \tilde{\Gamma}_0 f, \tilde{\Gamma}_0 g \rangle_{N_+} - i \langle B^* \tilde{\Gamma}_0 f, B^* \tilde{\Gamma}_0 g \rangle_{N_-},$$

so L_B^* is anti-dissipative if and only if $BB^* \geq I_{N_+}$. Assuming L_B is maximally dissipative, we obtain the Lagrange identities (3.1) and (3.2) by choosing $\Gamma = (B^* B - I_{N_-})^{1/2} \Gamma_0$, $\Gamma_* = (BB^* - I_{N_+})^{1/2} \Gamma_1$, $E = \overline{\text{Ran } \Gamma}$ and $E_* = \overline{\text{Ran } \Gamma_*}$.

Now, for $f = f_0 + Bf_{-i} + f_{-i} \in D(L_B)$, we have

$$\begin{aligned} S(\lambda)(B^* B - I_{N_-})^{1/2} \Gamma_0 f &= S(\lambda) \Gamma f \\ &= \Gamma_*(L_B^* - \lambda)^{-1} (L_B - \lambda)(f_0 + Bf_{-i} + f_{-i}) \\ &= (BB^* - I_{N_+})^{1/2} \Gamma_1 (f_0 + (L_B^* - \lambda)^{-1} ((i - \lambda) Bf_{-i} - (i + \lambda) f_{-i})) \\ &= \frac{1}{\sqrt{2}} (BB^* - I_{N_+})^{1/2} \Gamma_1 (L_B^* - \lambda)^{-1} ((i - \lambda) B\Gamma_0 f - (i + \lambda) \Gamma_0 f). \end{aligned}$$

Since $\Gamma_0 u \in N_-$ is arbitrary, we get for any $f_{-i} \in N_-$ that

$$(8.6) \quad S(\lambda)(B^* B - I_{N_-})^{1/2} f_{-i} = \frac{1}{\sqrt{2}} (BB^* - I_{N_+})^{1/2} \Gamma_1 ((L_B^* - \lambda)^{-1} ((i - \lambda) Bf_{-i} - (i + \lambda) f_{-i})).$$

From now on, assume $f \in \ker(L^* - \lambda)$. Then

$$(L - \lambda)f_0 + (i - \lambda)f_i - (i + \lambda)f_{-i} = 0$$

and so

$$\Gamma_1(L_B^* - \lambda)^{-1}((i - \lambda)f_i - (i + \lambda)f_{-i}) = 0.$$

This allows us to rewrite (8.6) as

$$(8.7) \quad S(\lambda)(B^*B - I_{N_-})^{1/2}f_{-i} = \frac{i - \lambda}{\sqrt{2}}(BB^* - I_{N_+})^{1/2}\Gamma_1(L_B^* - \lambda)^{-1}(Bf_{-i} - f_i).$$

Now let $\widetilde{M}(\lambda)$ be the M -function such that

$$\widetilde{M}(\lambda)(\widetilde{\Gamma}_1 - B^*\widetilde{\Gamma}_0)f = \widetilde{\Gamma}_0f \quad \text{for } f \in \ker(L^* - \lambda).$$

Then for such $f \in \ker(L^* - \lambda)$, we have

$$(8.8) \quad \widetilde{M}(\lambda)(f_{-i} - B^*f_i) = f_i.$$

Setting $w_+ = Bf_{-i} - f_i \in N_+$ and $h = (L_B^* - \lambda)^{-1}w_+$, we get

$$(L^* - \lambda) \left(h - \frac{w_+}{i - \lambda} \right) = 0.$$

This implies

$$\widetilde{M}(\lambda) \left(h_{-i} - B^*(h_i - \frac{w_+}{i - \lambda}) \right) = h_i - \frac{w_+}{i - \lambda}.$$

Since $h \in D(L_B^*)$, we have $h_{-i} = B^*h_i$, and therefore we get

$$h_i = (I + \widetilde{M}(\lambda)B^*)\frac{w_+}{i - \lambda}.$$

Inserting this in (8.7), we get

$$(8.9) \quad S(\lambda)(B^*B - I_{N_-})^{1/2}f_{-i} = (BB^* - I_{N_+})^{1/2}(I + \widetilde{M}(\lambda)B^*)(Bf_{-i} - f_i).$$

From (8.8), we have that $(I + \widetilde{M}(\lambda)B^*)f_i = \widetilde{M}(\lambda)f_{-i}$, so

$$\begin{aligned} S(\lambda)(B^*B - I_{N_-})^{1/2}f_{-i} &= (BB^* - I_{N_+})^{1/2}(B + \widetilde{M}(\lambda)B^*B - \widetilde{M}(\lambda))f_{-i} \\ &= \left[B + (BB^* - I_{N_+})^{1/2}\widetilde{M}(\lambda)(B^*B - I_{N_-})^{1/2} \right] (B^*B - I_{N_-})^{1/2}f_{-i}. \end{aligned}$$

Therefore, on N_- , we have

$$(8.10) \quad S(\lambda) = B + (BB^* - I_{N_+})^{1/2}\widetilde{M}(\lambda)(B^*B - I_{N_-})^{1/2}.$$

Remark 8.2. (1) *This gives a generalisation of Ryzhov's formula for the characteristic function to the case of different deficiency indices.*

(2) *If L is not symmetric, the general formula will be much more complicated, as seen e.g. in (4.11).*

(3) *Rewriting the expressions in terms of the M -function $\widetilde{M}_\infty(\lambda)$ gives an easier expression. Note that $\widetilde{M}_\infty(\lambda)$ is a contraction.*

(4) *Our choice of boundary triple here is not the 'expected' one for an underlying symmetric operator, therefore our M -function is not Nevanlinna. Instead it is contractive and L_B is dissipative iff B is contractive.*

(5) *If the deficiency indices coincide we are able to choose the Γ -operators to get Ryzhov's representation and have an M -function that is Nevanlinna.*

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