Time-Reversal Symmetry Breaking in Re-Based Superconductors

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To trace the origin of time-reversal symmetry breaking (TRSB) in Re-based superconductors, we performed comparative muon-spin rotation and relaxation (μSR) studies of superconducting noncentrosymmetric Re0.82 Nb0.18 (Tc = 8.8 K) and centrosymmetric Re (Tc = 2.7 K). In Re0.82 Nb0.18, the low-temperature superfluid density and the electronic specific heat evidence a fully gapped superconducting state, whose enhanced gap magnitude and specific-heat discontinuity suggest a moderately strong electron-phonon coupling. In both Re0.82 Nb0.18 and pure Re, the spontaneous magnetic fields revealed by zero-field μSR below Tc indicate time-reversal symmetry breaking and thus unconventional superconductivity. The concomitant occurrence of TRSB in centrosymmetric Re and noncentrosymmetric ReT (T = transition metal), yet its preservation in the isostructural noncentrosymmetric superconductors Mg10 Ir19 B16 and Nb0.5 Os0.5, strongly suggests that the local electronic structure of Re is crucial for understanding the TRSB superconducting state in Re and ReT. We discuss the superconducting order parameter symmetries that are compatible with the experimental observations.

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Time reversal and spatial inversion are two key symmetries that influence, at a fundamental level, the electron pairing in the superconducting state. On the one hand, a number of unconventional superconductors exhibit spontaneous time-reversal symmetry breaking (TRSB) on entering the superconducting state; on the other hand, the absence of inversion symmetry already above Tc leads to an antisymmetric spin-orbit coupling (SOC), lifting the degeneracy of the conduction-band electrons and potentially giving rise to a mixed-parity superconducting state [1,2]. Some noncentrosymmetric superconductors (NCSC), such as CePt3 Si [3], CeIrSi3 [4], Li3Pt3 B [5,6], and K2Cr3As3 [7,8], exhibit line nodes in the gap, while others, such as LaNiC2 [9] and (La, Y)2Cu3 [10], show multiple nodeless superconducting gaps. In addition, due to the strong influence of SOC, their upper critical field can greatly exceed the Pauli limit, as has been found in CePt3 Si [11] and very recently in (Ta, Nb)Rh2B2 [12].

In general, TRSB below Tc and a lack of spatial-inversion symmetry of the crystal structure are independent events. Yet, in a few cases, such as in LaNiC2 [13], La3Ir3 [14], and in particular, in the Re-based compounds Re7Zr [15], Re8Hf [16], Re8Ti [17], and Re24Ti8 [18], TRSB below Tc is concomitant with an existing lack of crystal inversion symmetry. Such an unusually frequent occurrence of TRSB among the superconducting ReT binary alloys (T = transition metal) is rather puzzling. Its persistence, independent of the particular transition metal, points to a key role played by Re. To test such a hypothesis, and to ascertain the possible relevance of the noncentrosymmetric structure to TRSB in Re-based NCSC, we proceeded with a twofold study. On one hand, we synthesized and investigated another Re-based NCSC, Re0.82 Nb0.18. On the other hand, we considered the pure Re metal, also a superconductor, but with a centrosymmetric structure.
A comparative study by means of muon-spin relaxation and rotation ($\mu$SR) allows us to address the question of TRSB in Re-containing compounds. The choice of $\mu$SR as the preferred technique for our study is justified by its key role in detecting TRSB in numerous unconventional superconductors [13–20] (later confirmed by the Kerr effect or bulk magnetization in the cases of Sr$_2$RuO$_4$, UPt$_3$, and LaNiC$_2$ [21–23]). We report systematic $\mu$SR studies of Re$_{0.82}$Nb$_{0.18}$ ($T_c = 8.8$ K) and Re ($T_c = 2.7$ K), whose bulk superconducting properties were characterized by magnetic, transport, and thermodynamic measurements. The $\mu$SR data show that spontaneous magnetic fields appear below the respective transition temperatures, thus implying that the superconducting states of both Re$_{0.82}$Nb$_{0.18}$ and Re show TRSB and have an unconventional nature. Since pure Re is centrosymmetric, this implies that the noncentrosymmetric structure is not a requirement for TRSB in these materials.

Polycrystalline Re$_{0.82}$Nb$_{0.18}$ samples were prepared by arc melting Re and Nb metals and the same Re powder was used for measurements on elementary Re. The x-ray powder diffraction measured using a Bruker D8 diffractometer, confirmed the $\alpha$-Mn structure of Re$_{0.82}$Nb$_{0.18}$ ($I43m$) and the hcp-Mg structure of Re (P6$_3$/mmc) [24–30]. Magnetic susceptibility, electrical resistivity, and specific-heat measurements were performed on a 7-T Quantum Design Magnetic Property Measurement System and a 9-T Physical Property Measurement System. The $\mu$SR measurements were carried out on both the MuSR instrument of the ISIS pulsed muon source (UK) [31], and the GPS and LTF spectrometers of the $\pi$M3 beam line at the Paul Scherrer Institut, Villigen, Switzerland.

The magnetic susceptibility was determined using both field-cooled (FC) and zero-field-cooled (ZFC) protocols. As shown in Figs. 1(a)–1(b), the splitting of the two curves is typical of type-II superconductors, with the ZFC-susceptibility indicating bulk superconductivity with $T_c = 8.8$ K for Re$_{0.82}$Nb$_{0.18}$ and 2.7 K for Re. The bulk superconductivity of Re$_{0.82}$Nb$_{0.18}$ was further confirmed by electrical resistivity and specific-heat data [24]. To perform transverse-field muon-spin rotation (TF-$\mu$SR) measurements of superconductors, the applied field should exceed the lower critical field $H_{c1}$, so that the additional field-broadening due to the flux-line lattice (FLL), can be determined from the decay of the $\mu$SR asymmetry. To determine $\mu_0H_{c1}$, the field-dependent magnetization $M(H)$ was measured at various temperatures below $T_c$, as shown in Fig. 1(c), for Re$_{0.82}$Nb$_{0.18}$ [for $M(H)$ data of Re, see Supplemental Material] [24]. The derived $\mu_0H_{c1}$ values are plotted in Fig. 1(d) as a function of temperature. The dashed lines are fits to $\mu_0H_{c1}(T) = \mu_0H_{c1}(0)[1 - (T/T_c)^2]$, which yield estimates of lower critical fields of 6.4(1) mT and 3.7(2) mT in Re$_{0.82}$Nb$_{0.18}$ and Re, respectively.

The TF-$\mu$SR measurements allowed us to explore the nature of superconductivity in Re$_{0.82}$Nb$_{0.18}$ at a microscopic level. The optimal field value for such experiments (above $H_{c1}$) was determined via preliminary field-dependent $\mu$SR measurements at 1.5 K [24]. Figure 2(a) shows two representative TF-$\mu$SR spectra collected above and below $T_c$ in an applied field of 15 mT. In the superconducting mixed state, the faster decay of muon-spin polarization reflects the inhomogeneous field distribution due to the FLL. The corresponding TF spectra are described by:

$$P_{TF} = P_s \cos(\gamma \mu B_s t + \phi)e^{-\sigma^2 t^2/2} + P_{bg} \cos(\gamma \mu B_{bg} t + \phi).$$

Here, $P_s$ and $P_{bg}$ represent the muon-spin polarization for muons implanted in the sample and sample holder, respectively, with the latter not undergoing any depolarization. $\gamma = 2\pi \times 135.53$ MHz/T is the muon gyromagnetic ratio, $B_s$ and $B_{bg}$ are the respective local fields sensed by implanted muons in the sample and sample holder, $\phi$ is the initial phase, and $\sigma$ is a Gaussian relaxation rate. In the superconducting state, the Gaussian relaxation rate includes contributions from both the FLL ($\sigma_{sc}$) and a temperature-independent relaxation due to nuclear moments ($\sigma_n$). Below $T_c$, $\sigma_{sc}$ can be extracted after subtracting $\sigma_n$ in quadrature.
Figure 2(b) shows the normalized superfluid density \( (\lambda(T)/\lambda(0))^{-2} \) vs the reduced temperature \( T/T_c \) for Re\(_{0.82}\)Nb\(_{0.18}\). Clearly the temperature dependence of the superfluid density is highly consistent between PSI and ISIS, and well described by an s-wave model with a single gap of 1.61(1) meV. By using the 15-mT data, the resulting \( \lambda(0) \) of 357(3) nm is comparable with 352(3) nm, the value calculated from \( \mu_0 H_{c1} \) [24]. The superconducting gap is similar to that of other ReT superconductors, e.g., Re\(_{0.75}\)Zr (1.21 meV) [15,34], Re\(_{0.75}\)Ti\(_x\) (1.08 meV) [18], Re\(_{0.75}\)Ti (0.95 meV) [17], and Re\(_{0.75}\)Hf (1.10 meV) [16,35,36], (see also Table SI in the Supplemental Material) [24]. Also, the \( 2\Delta/k_B T \) values of these compounds [e.g., 4.26 for Re\(_{0.82}\)Nb\(_{0.18}\)] are higher than 3.53, the value expected for weakly-coupled BCS superconductors, thus indicating a moderately strong electron-phonon coupling in these materials. The superconducting parameters of all these \( \alpha\)-Mn-type ReT NCSC are summarized in Table SI in the Supplemental Material [24].

A detailed analysis of the zero-field specific-heat data provides further insight into the superconducting properties of Re\(_{0.82}\)Nb\(_{0.18}\). The electronic specific heat \( C_e/T \) was obtained by subtracting the phonon contribution from the experimental data [24]. The derived \( C_e/T \) was then divided by the normal-state electronic specific heat coefficient, as shown in the main panel as a function of \( T/T_c \). The solid line in Fig. 2(c) represents a fit with \( \gamma_n = 4.4 \) mJ mol\(^{-1}\) K\(^{-2}\) and a single isotropic gap \( \Delta(0) = 1.52(2) \) meV. This reproduces the experimental data very well, while being consistent with the TF-\( \mu \)SR [see Fig. 2(b)] and previously reported values [37,38]. Note also that this value is between the two values found from a two-gap analysis [39]. The specific-heat jump at \( T_c \) was found to be \( \Delta C/\gamma_n T_c \sim 1.94 \), i.e., larger than the conventional BCS value of 1.43, again indicating a moderately enhanced electron-phonon coupling in Re\(_{0.82}\)Nb\(_{0.18}\).

The key goal of the present Letter is to probe a possible TRSB in Re\(_{0.82}\)Nb\(_{0.18}\) and in pure Re. To this aim, we performed detailed zero-field muon-spin relaxation (ZF-\( \mu \)SR) measurements. Normally, in the absence of external fields, the onset of superconductivity does not imply changes in the ZF-\( \mu \)SR relaxation rate. However, in the presence of TRSB, the onset of a tiny spontaneous polarization or currents gives rise to associated (weak) magnetic fields, readily detected by ZF-\( \mu \)SR as an increase in the \( \mu \)SR relaxation rate. Given the tiny size of such effects, we measured the ZF-\( \mu \)SR, both above \( T_c \) and inside the superconducting phase. Representative ZF-\( \mu \)SR spectra collected above and below \( T_c \) for Re\(_{0.82}\)Nb\(_{0.18}\) and Re show measurable differences [see Figs. 3(a)–3(b)]. To exclude the presence of stray magnetic fields, all magnets were quenched before the measurements and an active compensation system was used. In nonmagnetic materials in the absence of applied fields, the relaxation is determined primarily by the randomly oriented nuclear dipole.
shown in the insets of Fig. 3, despite the different below
of Re
nuclear magnetic moment for Re
ISIS in (a) were obtained by fitting the ZF data with the corre-
the respective nuclear moments
indicates a linear behavior. The nuclear moment was estimated from
field, are also shown. The solid lines are fits using Eq.(2). Insets
relaxation function[40,41]. In our case, the ZF-
moments, normally described by a Gaussian Kubo-Toyabe
additional μSR data sets collected at 1.5 K in a 15-mT longitudinal
field, are well described by a combined Lorentzian and Gaussian
TRSB to the specific transition-metal element
have a negligible effect on TRSB. The insensitivity of
spontaneous magnetic fields in both samples, indicating
that the relevant fields are static on the time scale of the
muon lifetime.

To date, most α-Mn-type ReT-NCSC have been found
to exhibit TRSB in the superconducting state [15–18]. Our
results show that Re0.82Nb0.18 is not just another member of the ReT-NCSC family, but one with the most
distinct TRSB in the superconducting state (i.e., with the
highest σint, which represents the change of muon
relaxation rate between the normal and superconducting
states (see details in Table SI in Supplemental Material)
[24]. This is clearly depicted in Fig. 3(c), where we plot
the estimated internal field Bint as a function of the
nuclear magnetic moment μn for the ReT-NCSC. It can
be seen that, Bint (∝ σint) scales linearly with μn, reaching
0.038 mT for Re0.82Nb0.18. At the other extreme, the line
crosses the horizontal axis at μn ∼ 2.7μN, where σint drops
below the resolution of the μSR technique (0.01 mT).
This is exactly the case for Re3W, whose σint turned out
to be negligible [42].

Having detected TRSB in Re-based superconductors still
leaves open the most intriguing question: what is its key
ingredient? In ReT, the replacement of heavy 5d elements,
such as Hf, with lighter 3d elements, such as Ti, appears to
have a negligible effect on TRSB. The insensitivity of
TRSB to the specific transition-metal element T suggests
that a substitution at the T-sites does not significantly
influence it. This is confirmed by the persistence of TRSB
in elemental Re. In addition, this indicates also that a lack
of inversion symmetry is inessential. Finally, there is no
TRSB in the superconducting states of Mg2Ir1B16 [43]
and Nb5O8.5 [44], two NCSCs isostuctural to ReT and
with similar SOC strengths. The above considerations
strongly suggest that it is the local electronic structure
of Re that is crucial for understanding the TRSB in the
superconducting states of Re and ReT. To reinforce the
above conclusion, one could study other Re-free materials
with the α-Mn-type structure. TaOs, with a bulk Tc of
other ReT NCSC [15–18]. At the same time, the Lorentzian
relaxation rate Λ(T) remains mostly constant in the studied
temperature range, with typical values of 0.007 and
0.005 μS−1 for Re0.82Nb0.18 and Re [24], respectively,
indicating that fast-fluctuation effects are absent in these
systems. The small, yet measurable increases of σ(T)
below Tc, detected from measurements at both facilities,
reflect the onset of spontaneous magnetic fields, and thus,
they are signatures of TRSB in the superconducting phases
of both Re0.82Nb0.18 and Re. Further refinements performed
by fixing the Λ values gave similarly robust features in
σ(T) [24]. To rule out the possibility of a defect- or
impurity-induced relaxation at low temperatures, we per-
formed auxiliary longitudinal-field μSR measurements at
1.5 K. As shown in Figs. 3(a)–3(b), a small field of 15 mT
is sufficient to decouple the muon spins from the weak
spontaneous magnetic fields in both samples, indicating
that the relevant fields are static on the time scale of the
muon lifetime.

![Graph and text](image-url)

FIG. 3. Representative zero-field μSR spectra for (a) Re0.82Nb0.18 and (b) pure Re metal in the superconducting and normal states. Additional μSR data sets collected at 1.5 K in a 15-mT longitudinal field, are also shown. The solid lines are fits using Eq. (2). Insets show the T-dependence of the relaxation rate σ. The results from ISIS in (a) were obtained by fitting the ZF data with the corres-
ponding expression in Ref. [15]. (c) Calculated internal field vs nuclear magnetic moment for ReT superconductors. For the pure Re (diamond), the value was obtained by an extrapolation to 0 K. The reported data are from Refs. [15–18,34–36,42]. The dashed-line indicates a linear behavior. The nuclear moment was estimated from the respective nuclear moments μn,Re, μn,T and chemical fractions fRe and fT of Re and T using μn = \sqrt{fReμn,Re + fTμn,T}.

moments, normally described by a Gaussian Kubo-Toyabe
relaxation function [40,41]. In our case, the ZF-μSR spectra
are well described by a combined Lorentzian and Gaussian
Kubo-Toyabe relaxation function:

\[
P_{\text{CKT}} = P_s \left[ \frac{1}{3} + \frac{2}{3} (1 - \sigma^2)^2 - \Lambda t e^{-[\sigma^2/2] - \Lambda t} \right] + P_{\text{bg}}. \tag{2}\]

Here, \( P_s \) and \( P_{\text{bg}} \) are the same as in the TF-μSR spectra. As shown in the insets of Fig. 3, despite the different Tc values of Re0.82Nb0.18 and Re, their \( \sigma(T) \) curves exhibit a small yet distinct increase below Tc, similar to that found also in
2.07 K and the required crystal structure, represents a good example [45].

We now discuss the possible symmetries of the superconducting order parameter. In the limit of weak SOC, TRSB can be achieved via nonunitary triplet pairing, as, e.g., in LaNiC$_2$ [46] and LaNiGa$_2$ [47,48]. More generally, the relationship between TRSB and triplet pairing is quite complex. For example, the admixed triplet component in so-called “s-wave” NCSC (those whose superconducting instabilities do not break crystal point-group symmetries), such as Li$_3$Pt$_3$B, does not show TRSB [5]. Conversely, there are TRSB states not involving triplet pairing, e.g., the $s + id$ singlet state proposed for some iron-based superconductors [49]. Apart from the weak-SOC nonunitary triplet pairing scenario mentioned above [46,47], the essential requirement for TRSB occurrence is that the point group of the crystal has irreducible representations with dimension $D > 1$.

The point groups $T_d$ and $D_{4h}$ relevant to ReT and Re, respectively, have several irreducible representations with $D = 2$ or 3. Therefore, they can support TRSB states with singlet-, triplet-, or in the case of ReT, admixed pairing, independent of SOC strength. In what follows, we will assume strong SOC. The full symmetry analysis and plots of the possible order parameters can be found in the Supplemental Material [24]. For ReT there are a number of possible TRSB states, with some examples of pairing functions being given in Ref. [15]. However, all of the possible states have symmetry-constrained point or line nodes, inconsistent with the experimental observations. In view of this, in some systems, it has been proposed that a full gap may be obtained through a Loop-Josephson-Current (LJC) state built on site, intraorbital, singlet pairing. Although it has been shown that the crystal symmetry of ReT is compatible with this scenario [50], the energetics that would drive such a state, if realized, and why it would occur only in systems with Re and not other elements, remain unclear.

The symmetry analysis of pure Re contrasts strongly that of ReT. First, Re is centrosymmetric, implying that the superconducting instability can only take place in either purely-singlet or purely-triplet channels (irrespective of the strength of SOC). Second, due to only two distinct symmetry-related sites per unit cell, a LJC state here is not the most natural one [50]. Third, the crystallographic space group is nonsymmorphic, which in principle allows superconducting instabilities that break screw-axis or glide-plane symmetries. Ignoring the ones that break those symmetries, we find two possible TRSB states, one in the singlet channel with a line node at $k_z = 0$, and another in the triplet channel, with two point nodes on the $k_z$ axis (see Supplemental Material [24]). The Fermi surface of Re has five sheets, including an electron sheet centered on the Γ point and open along the $k_z$ axis, and three hole sheets that are closed, centered on the $L$ point, and not intersecting the $k_z = 0$ plane [51]. These would be compatible with a full gap for the triplet and singlet TRSB states, respectively.

In summary, we performed comparative $\mu$SR studies of the noncentrosymmetric Re$_{0.82}$Nb$_{0.18}$ and centrosymmetric Re superconductors. Bulk superconductivity with $T_c = 8.8$ K (Re$_{0.82}$Nb$_{0.18}$) and 2.7 K (Re) was characterized by magnetic and transport properties. Both the superfluid density and the zero-field specific-heat data reveal a single-gap nodeless superconductivity in Re$_{0.82}$Nb$_{0.18}$. The spontaneous fields appearing below $T_c$, which increase with decreasing temperature, provide strong evidence that the superconducting states of both noncentrosymmetric Re$_{0.82}$Nb$_{0.18}$ and centrosymmetric Re show TRSB and are unconventional. Comparisons with other Re-free α-Mn-type superconductors suggest that in the ReT family, the TRSB is crucially related to the presence of Re, a key idea for understanding the peculiar behavior of ReT superconductors. We have considered the possible symmetries of the order parameter in these systems and their compatibility with the observed fully gapped spectrum. Further theoretical and experimental work on Re is required to clarify the open issues.

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[24] See the Supplemental Material at http://link.aps.org/supplemental/10.1103/PhysRevLett.121.257002 for details on the measurements of crystal structure, heat capacity, electrical resistivity, and critical field, as well as for the data analysis and the symmetry-allowed TRSB order parameters, which includes Refs. [25–30].


